Inverse and enhanced magnetocaloric properties of HoCrO₃

Ganesh Kotnana, Dwipak Prasad Sahu and S. Narayana Jammalamadaka*

Magnetic Materials and Device Physics Laboratory, Department of Physics, Indian Institute

of Technology Hyderabad, Hyderabad, India – 502 285.

*Corresponding author: surya@iith.ac.in

Abstract:

We report on the magnetic and magnetocaloric properties of 50% Fe³⁺ doped polycrystalline

HoCrO₃ compounds in the vicinity of magnetic transitions. Due to complex magnetic

interactions, we do see multiple transitions in χ vs. T graph pertinent to HoCrO₃ and

HoCr_{0.5}Fe_{0.5}O₃ compounds related with Ho³⁺ ordering, Cr³⁺ ordering and spin – reorientation

(SR). Due to re-orientation of spins, metamagnetic transitions are evident in low field regime

of M vs. H. Quantified values of -ΔS_M around Ho³⁺ ordering infer that indeed Fe³⁺

substitution helped in enhancing magnetocaloric effect of HoCrO₃ compound. Such an

enhanced -ΔS_M values are ascribed to increase in canting of Fe/Cr spins. Present results

would be helpful in developing magnetic refrigerant materials for space applications

particularly below 30 K.

Keywords: Magnetocaloric effect, magnetization crossover, metamagnetic transition, spin –

re-orientation

1

Introduction:

In recent years, refrigeration based on magnetic materials particularly using magnetocaloric effect (MCE) phenomenon has gained much attention in view of its application in future magnetic refrigeration technology due to its higher energy efficiency and environmental safety over the conventional gas compression refrigeration¹⁻³. MCE is a magneto-thermal phenomenon, in which the temperature of a magnetic material changes when it is exposed to a magnetic field adiabatically⁴. A magnetic materials with a large change of magnetic entropy $\Delta S_M(T,H)$ or adiabatic temperature (ΔT) and a sufficiently large value of refrigeration capacity (RC) are promising candidates for magnetic refrigeration⁵. A plethora of compounds have been investigated to study the MCE in alloys like Ni_{52.6}Mn_{23.1}Ga_{24.3}⁶, Gd₅(Si₂Ge₂)⁷, MnFeP_{0.45}As_{0.55}⁸ as well as in oxide materials such as Tb_{0.67}R_{0.33}MnO₃ (R = Dy, Y and Ho)⁹ and La_{1-x}Ca_xMnO₃¹⁰.

As for as oxide materials are concerned, distorted perovskite family of canted antiferromagnets, orthoferrites (RFeO₃)¹¹ and orthochromites (RCrO₃)¹² have attained much interest due to their unique physical properties and complex spin structure between the rareearth (R) and the transition metal (TM) ions. RFeO₃ and RCrO₃ compounds have been explored for their MCE properties and suitability for magnetic refrigeration applications¹³⁻¹⁵. For example, a large MCE value of Δ S ~ 8.4 J/kg-K and RC ~ 217 J/kg at 15 K and 4 T has been observed in DyCrO₃ due to ordering of Dy³⁺ at 2.3 K¹⁴. Yin *et.al*, reported large MCE value of Δ S ~ 10.5 J/kg-K at 4 T around Dy³⁺ ordering (14 K) in DyFe_{0.5}Cr_{0.5}O₃ compound¹⁵. Yet in another compound, HoCrO₃, large value of MCE has been reported and the values are Δ S ~7.2 J/kg-K and RC ~ 408 J/kg at 20 K and at 7 T¹⁶. Although several reports have shown MCE on various magnetic materials like Dy_{1-x}Sc_xNi₂¹⁷, R_{1-x}La_xNi₂ (R = Tb, Dy, Ho)¹⁸, Dy₂Cu₂Cd¹⁹, TmAgAl²⁰, TmZn²¹, Er₂Cu₂In²², Eu₄PdMg²³, etc... as this field is highly

demanding numerous new materials needs to be explored for various applications. Particularly, much attention is required to search for new materials with giant MCE below 30 K for utilization in space applications^{1, 24}.

In this context, HoCrO₃ is a compound which shows large value of ΔS (~7.2 J/kg-K) below 30 K. It is evident from the above discussion that substitution of Fe³⁺ in DyCrO₃ for Cr³⁺ has enhanced $-\Delta S$ values around Dv^{3+} ordering 15. Hence, we believe that Fe^{3+} would play a crucial role in enhancing MCE in our present HoCr_{0.5}Fe_{0.5}O₃ compound around Ho³⁺ ordering. In addition, it also has been believed from neutron diffraction data that up on adding 50% of Fe to Cr, spin reorientation (SR) transition would be evident²⁵ due to an interaction between Ho³⁺ and Fe³⁺ moments. Such a SR transition would be associated with a change in magnetization direction from $G_xA_vF_z$ (Γ_4) to $F_xC_vG_z$ (Γ_2), where G_x , A_v , and F_z stand for spin component along x, y, and z axis in terms of Bertaut's notation, respectively²⁶, which can lead to large change in -ΔS_M. Calculation of entropy change around SR transition is also aim of our present paper. Apart from the SR and Ho³⁺ ordering, indeed there exists a transition due to Cr³⁺, which may also give large change in MCE. Hence, in the present work, we explored MCE properties of HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds around Ho³⁺, Cr3+ and SR transition temperatures. In addition, we also studied the nature of magnetic transition in aforesaid compounds through Arrott plots. Sailent features of present manuscript are (a) up on adding 50% of Fe³⁺ to Cr³⁺, enhanced value of MCE is evident in HoCr_{0.5}Fe_{0.5}O₃ compound (b) magnetization crossover is evident in HoCrO₃ compound and metamagnetic transition prevails around Ho³⁺ ordering in both HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds (c) Spin reorientation (SR) transition is apparent in HoCr_{0.5}Fe_{0.5}O₃ compound around 150 K.

Experimental Details:

Polycrystalline HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds were prepared by conventional solid state reaction method using High purity oxide powders of Ho₂O₃, Fe₂O₃, and Cr₂O₃ (purity > 99.9%) (Sigma-Aldrich chemicals India) as starting raw materials. The temperature (T) dependent magnetization (M) (M vs. T) and magnetic field (H) dependent magnetization (M vs. H) measurements were performed using a Quantum Design magnetic property measurement system (MPMS) in the temperature range of 5 – 300 K. Magnetization isotherms (M vs. H) were measured at different temperature ranges and up to a maximum magnetic field of 6 T. Before we perform each measurement of M vs. H, the sample was warmed upto a temperature greater than ordering temperature of transition metal to remove the magnetic history.

Results and Discussion:

The phase purity of HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds were confirmed²⁷ at room temperature using powder x - ray diffraction (XRD) (PANalytical X-ray diffractometer) with Cu - K_{α} radiation (λ = 1.5406 Å) and with a step size of 0.017° in the wide range of the Bragg angles 2 θ (20° - 80°) as shown in Fig. 1. From the Fig.1, it is observed that the indexed reflections are allowed for a compound with GdFeO₃ type distorted perovskite structure with a space group of *Pbnm*. We do not see any impurity phase within the detectable limits of the XRD. The structural parameters obtained from the Rietveld refinement has been published elsewhere²⁷. The increase in the lattice parameter is evident from the refinement with increasing Cr³⁺ content, which is consistent with the fact that ionic radius of Cr³⁺ (0.615 Å) is smaller than that of Fe³⁺ (0.645 Å).

Fig. 2(a) shows the temperature (T) vs. susceptibility (χ) graph in zero field cooled (ZFC) conditions. It is evident from the figure that the Cr^{3+} moments ordered antiferromagnetically (AFM) at 142 K. Such AFM is believed have canting nature with AFM axis at 142 K (T_N). On the other hand, transition due to Ho^{3+} moments is evident at 8 K (T_N^{Ho})²⁸. Fig. 2(b) shows χvs . T graph for $HoCr_{0.5}Fe_{0.5}O_3$ compound. In 50% Fe substituted compound, in addition to AFM and Ho^{3+} (13 K) ordering, we also could able to realize a transition due to SR^{25} . The origin of such SR is due to change in magnetization from $G_xA_yF_z$ (Γ_4) to $F_xC_yG_z$ (Γ_2), where G_x , A_y , and F_z stand for spin component along x, y, and z axis in terms of Bertaut's notation, respectively. Such a SR transition is evident around 150 K (T_{SR}) and we also ascribe such SR in $HoCr_{0.5}Fe_{0.5}O_3$ compounds to complex exchange interaction between the Fe^{3+} and the Ho^{3+} ions²⁹. Indeed, there is a shift for Ho^{3+} ordering from 8 – 13 K and T_N from 142 – 273 K up on adding 50% Fe to Cr. The obtained T_N and T_{SR} values are in good agreement with the literature^{25, 28}.

As there are multiple transitions in both the compounds, we do expect a large change in entropy around these transitions. Hence, in order to quantify change in entropy we performed the M vs. H measurements in the vicinity of the magnetic transitions for both the HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃compounds. Fig. 3(a)-(d) shows first quadrant magnetization isotherms of HoCrO₃ (at 2 – 10 K and 136 – 146 K with an interval of 2 K) and HoCr_{0.5}Fe_{0.5}O₃ (at 5 - 20K with an interval of 3 K and 120 – 200 K with an interval of 10 K) compounds with a maximum magnetic field of 6 T. A field induced metamagnetic transition is evident in the low magnetic field region for all M vs. H graphs pertinent to HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds as shown in Fig. 3(a) and 3(b). Now we define the critical field (H_C) as field below which the material shows AFM behaviour and exhibits ferromagnetic (FM) behaviour above this field. Such a field induced transition in low field region can be ascribed to the onset of a first order metamagnetic transition (from original AFM state to FM state).

Evidenced metamagnetic transition may be due to the magnetization reversal of those ions whose magnetic moments are directed opposite to the applied magnetic field and the similar behaviour has been observed in the single crystals of DyCrO₃ as the rotation of Dy³⁺ and Cr³⁺ moments³⁰. From Fig.3 (a), it is evident that there is a magnetization crossover (which is shown as an inset of Fig. 3(a)) around the AFM \rightarrow FM transition, which can be ascribed as a result of the competition between the Zeeman energy due to the applied external field and the strong magneto-crystalline anisotropy energy around the metamagnetic transition³¹.

H_C values are obtained from the first derivative of M vs. H curve in low field region as shown in inset of Fig. 3 (b). The variation of H_C with temperature for HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ is shown in Fig 4(a) and 4(b). In both the compounds initially H_C increases up to a transition pertinent to Ho³⁺ ions and above which it diminishes. The increase in H_C up to Ho³⁺ ordering temperature can be explained as follows. For a perfect antiferromagnetic (AFM) system, below transition, in principle the H_C should increase due to the strong AFM nature. However, we do see a decreasing trend in our sample due to the presence of both ferromagnetic (FM) and AFM phases. Neutron diffraction study on HoFeO₃ compound³² has been revealed that the Ho³⁺ moments order in a distorted antiferromagnetic (AFM) configuration and each Ho³⁺ moment consists a magnetic moment with a magnitude of 7.5 μ_B in (001) plane and makes an angle with respect to [010] direction so as to produce a net ferromagnetic moment of 3.4 µB per HoFeO₃ molecule parallel to [100] at 1.25 K. Hence, the observed behavior of H_C can be attributed to the temperature dependent competition between AFM and FM phases. Fig. 4 (c) depicts the variation of magnetization (M) with the applied magnetic field (H) recorded at 5 K for both HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds. Indeed, there is an enhanced magnetization in HoCr_{0.5}Fe_{0.5}O₃ compound compared with HoCrO₃ due to larger magnetic moment of Fe³⁺ ions compared to that of Cr³⁺ ions. The increase in FM component in

HoCr_{0.5}Fe_{0.5}O₃ compound might be a reason why we do see lower value of H_C in comparison with HoCrO₃.

Fig. 3 (c) depicts the M vs. H graphs around AFM transition in HoCrO₃ compound. It is evident that in low field region, indeed there exists a curvature due to a weak ferromagnetic component. However, such transition is very weak. Above 142 K, we do not see any curvature for M vs. H graphs at low field region as a result of paramagnetic phase. We also tried to map the variation of magnetization with magnetic field around SR transition in HoCr_{0.5}Fe_{0.5}O₃ compound. Fig. 3 (d) indicates M vs. H loops measured for HoCr_{0.5}Fe_{0.5}O₃ compound during spin reorientation (SR) transition.

As there are multiple transitions in both the compounds, we do expect a large change in entropy around these transitions. Hence, we calculated magnetic entropy change $(-\Delta S_M)$ from isothermal magnetization curves using well known Maxwell's equation as follows

$$\Delta S_M(T,H) = \int_0^{H_{\text{max}}} \left(\frac{\partial M}{\partial T}\right)_H dH - - - - (1),$$

where $H_{\rm max}$ is the maximum value of external applied field. From the above equation, it can be noticed that the value of $-\Delta S_{\rm M}$ depends on both values of magnetization (M) and $\left(\frac{\partial M}{\partial T}\right)_{\!\!H}$.

The larger values of $-\Delta S_M$ can be obtained when the values of M and $\left(\frac{\partial M}{\partial T}\right)_H$ are large in the magnitude³³. Quantified values of $-\Delta S_M$ and its temperature variation is shown in Fig. 5 (a) - (d). Here, Fig. 5(a) and 5(b) depicts the $-\Delta S_M$ vs. T around Ho³⁺ ordering for HoCrO₃ ($-\Delta S_M \sim 6.99$ J/kg.K at 9 K) and HoCr_{0.5}Fe_{0.5}O₃ ($-\Delta S_M \sim 8.18$ J/kg.K at 12.5 K) compounds respectively. Large values of $-\Delta S_M$ in both the compounds can be ascribed to evidenced metamagnetic transition and the Ho³⁺ ordering. Reason for the enhancement in the value of $-\Delta S_M$ in HoCr_{0.5}Fe_{0.5}O₃ compound compared to that of the HoCrO₃ compound can be

explained as follows. Essentially, HoCrO₃ is distorted type perovskite material and crystallizes in orthorhombic structure with a space group of Pbnm³⁴. For an ideal perovskite Cr-O1-Cr bond angle must be 180°. However, due to the distortion as a result of tilting of CrO₆ octahedra in HoCrO₃, there would be a change in Cr-O1-Cr bond angle (146.2°)²⁷, which is distinctly different from an ideal perovskite. This distortion decreases the orbital overlap and leads to non collinear antiferromagnetic structure of Cr3+ ions with a weak ferromagnetic component. Obtained bond angle value from Rietveld refinement pertinent to Fe(Cr)-O1-Fe(Cr) is smaller for HoCr_{0.5}Fe_{0.5}O₃ (145.3°) in comparison with the HoCrO₃ (146.2°) compound. This indicates an increase in canting of spins and hence an enhancement in the $-\Delta S_M$ value for HoCr_{0.5}Fe_{0.5}O₃ compound in comparison with HoCrO₃. The increase in canting angle of Fe/Cr sublattice with AFM axis for HoCr_{0.5}Fe_{0.5}O₃ has been evident from the neutron diffraction data on HoFe_{1-x}Cr_xO₃ solid solutions²⁵. The values of MCE obtained for HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds in the vicinity of Ho³⁺ ordering temperature are large when compared to DyMnO3 35 (- $\Delta S_M \sim 6.8$ J/kg.K at 7T and 10 K) , SmFe0.5Mn0.5O3 36 (5.6 J/kg.K at 18 K and with 7 T) and HoCr_{0.7}Fe_{0.3}O₃³⁷ (6.83 J/kg.K at 20 K and with 7 T). The MCE values are found to be smaller compared to that of DyCrO₃¹⁴ and DyFe_{0.5}Cr_{0.5}O₃¹⁵, which could be due to larger magnetic hysteresis exists in the present HoCr_{0.5}Fe_{0.5}O₃²⁷ compound, which leads to more energy loss in the thermal process and small values of MCE. However, the $-\Delta S_M$ obtained in the present study are smaller compared to giant MCE material such as TmZn²¹, ErMn₂Si₂³⁸, Tm₂Cu₂In²², HoPdIn³⁹ and RE₂Cu₂O₅ (RE = Dy and Ho)⁴⁰ in the low temperature region.

From the Fig. 5(a), it is evident that $-\Delta S_M$ values shows both positive and negative values as though there exists re-entrant magnetocaloric effect. Such an unusual nature of re-entrant inverse MCE in HoCrO₃ compound has been not observed earlier¹⁶. The reason for such an intriguing phenomenon may be due to the subtle changes in relative orientations of Ho³⁺ and

 Cr^{3+} moments occurs with a variation of temperature and/or magnetic field⁴¹. Earlier inverse magnetocaloric effect has been observed around SR transition in single crystals of HoFeO₃⁴². We conceive from our observation that magnetization crossover and complex interaction between FM and AFM phases might be a reason for re-entrant MCE behaviour for the present compound. From Fig. 5 (c) and 5(d), observed $-\Delta S_M$ value of 1.05 J/kg.K (around 141 K for HoCrO₃) and 0.61 J/kg.K (around 125 K for HoCr_{0.5}Fe_{0.5}O₃) at 6 T can be ascribed to the Cr^{3+} ordering and the spin-reorientation transition respectively.

It is believed that the magnitude of the magnetic entropy change at a particular magnetic phase transition and its dependence on temperature and magnetic field strongly depend upon the nature of the corresponding phase transition⁴³, so that it is essential to determine the nature of magnetic phase transitions in the compounds. To do this, the first quadrant magnetization isotherms were plotted in the form of Arrott plots⁴⁴ which can be deduced from Ginzburg-Landau theory in the close vicinity of the magnetic transition. The thermodynamic potential with Ginzburg-Landau type expansion which includes the magnetostatic field energy (*MH*) near the magnetic transition is as follows:

$$G(T,M) = G_0 + \alpha M^2 + \beta M^4 + \dots - MH$$
 ... (2)

where α , β are the Landau coefficients dependent on temperature. In the equilibrium condition, $\frac{\partial G}{\partial M} = 0$, the equation (2) reduces to $H/_{M} = \beta M^2 + \alpha$. According to Banerjee's criterion⁴⁵, a negative slope of $H/_{M}$ vs. M^2 at some point indicates the first order magnetic transition. On the other hand a positive slope of $H/_{M}$ vs. M^2 indicates the second order magnetic transition. Fig. 6(a)-(d) shows the Arrott plots corresponding to magnetization isotherms shown in Fig. 3(a)-(d) respectively. The negative slope of $H/_{M}$ vs. M^2 for the Fig. 6(a) - (c) corresponding to $HoCrO_3$ and $HoCrO_3$ compounds indicates the first order

magnetic phase transition, while the positive slope of $^{H}/_{M}$ vs. H for the Fig. 6(d) corresponding to H HoCr_{0.5}Fe_{0.5}O₃ indicates the second order magnetic phase transition which is in agreement with the literature⁴⁶.

Conclusions:

Due to complex magnetic interactions, we do see multiple transitions in χ vs. T graph of HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds. M vs. H graph infer that the magnetization crossover is evident around Ho³⁺ ordering in HoCrO₃. On the other hand, metamagnetic transitions are prevailed in the low field regime of M vs. H graph around Ho³⁺ ordering. Indeed, there is an enhancement in the saturation magnetization up on adding Fe³⁺ to Cr³⁺ due to large magnetic moment of Fe. Re-entrant magnetocaloric effect is evident in HoCrO₃, which may be due to magnetization crossover and complex interactions. We trust that the large values of MCE obtained for HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds makes them promising candidates for magnetic refrigeration applications below 30 K.

Acknowledgements

We would like to acknowledge Indian Institute of Technology, Hyderabad and Department of Science and Technology (DST) (Project #SR/FTP/PS-190/2012) for the financial support. We are grateful to Dr. Alok Banerjee, and Mr. Kranthi Kumar, UGC-CSR, Indore for carrying out the magnetization measurements. We also grateful to UGC-DAE Consortium (Project # CSR-IC/CRS-162/2015-16/19) for the financial support and their extended facilities.

References:

[1] Gschneidner Jr K. A, Pecharsky V. K and Tsokol A. O, Rep. Prog. Phys. 68, 1479 (2005).

- [2] Gutfleisch O, Willard M A, Bruck E, Chen C H, Sankar S G and Liu J P, *Adv. Mater.* 23, 821 (2011).
- [3] Shen B G, Sun J R, Hu F X, Zhang H W and Cheng Z H, Adv. Mater. 21, 4545 (2009).
- [4] N. S. Bingham, H. Wang, F. Qin, H. X. Peng, J. F. Sun, V. Franco, H. Srikanth and M. H. Phan, Appl. Phys. Lett. 101, 102407 (2012).
- [5] Ling-Wei Li, Chinese Phys. B 25, 037502 (2016).
- [6] Feng-xia Hu, Bao-gen Shen, Ji-rong Sun, and Guang-heng Wu, *Phys. Rev. B* **64**, 132412 (2001).
- [7] V. K. Pecharsky and K. A. Gschneidner Jr, *Phys. Rev. Lett.* **78**, 4494 (1997).
- [8] O. Tegus, E. Bruck, K. H. J. Buschow and F. R. de Boer, *Nature* **415**, 150-152 (2002).
- [9] M. Staruch, L. Kuna, A. McDannald and M. Jain, J. Magn. Magn. Mater. 377, 117–120 (2015).
- [10] Z. B. Guo, Y. W. Du, J. S. Zhu, H. Huang, W. P. Ding, and D. Feng, *Phys. Rev. Lett.* 78, 1142 (1997).
- [11] R. L. White, J. Appl. Phys. 40, 1061 (1969).
- [12] B. Rajeswaran, D. I. Khomskii, A. K. Zvezdin, C. N. R. Rao, and A. Sundaresan, *Phys. Rev. B* 86, 214409 (2012).
- [13] M. Shao, S. Cao, Y. Wang, S. Yuan, B. Kang and J. Zhang, Solid State Commun. 152, 947 (2012).
- [14] A. McDannald, L. Kuna and M. Jain, J. Appl. Phys. 114, 113904 (2013).

- [15] L. H. Yin, J. Yang, R. R. Zhang, J. M. Dai, W. H. Song and Y. P. Sun, Appl. Phys. Lett. 104, 032904 (2014).
- [16] S. Yin and M. Jain, J. Appl. Phys. **120**, 043906 (2016).
- [17] J. 'Cwik, T. Palewski, K. Nenkov, J. Lyubina, O. Gutfleisch, J. Klamut, J. Alloys Compd. 506, 626-630 (2010).
- [18] J. 'Cwik, Y. Koshkid'ko, I. Tereshina, N. Kolchugina, K. Nenkov, A. Hackemer, J. Lyubina, T. Palewski, G. S. Burkhanov, M. Miller, *J. Alloys Compd.* **649**, 417-425 (2015).
- [19] Scientific Reports 6 (2016) 34192
- [20] Yikun Zhang, Long Hou, Zhongming Ren, Xi Li, Gerhard Wilde, *J. Alloys Compd.* **656**, 635-639 (2016).
- [21] Lingwei Li, Ye Yuan, Yikun Zhang, Takahiro Namiki, Katsuhiko Nishimura, Rainer Pöttgen, and Shengqiang Zhou, *Appl. Phys. Lett.* **107**, 132401 (2015).
- [22] Yikun Zhang, Yang Yang, Xiao Xu, Long Hou, Zhongming Ren, Xi Li, and Gerhard Wilde, *J. Phys. D: Appl. Phys.* **49**, 145002 (2016).
- [23] D. H. Ryan, Anaëlle Legros, Oliver Niehaus, Rainer Pöttgen, J. M. Cadogan, and R. Flacau, *J. Appl. Phys.* **117**, 17D108 (2015).
- [24] N. A. de Oliveira, P. J. von Ranke, *Physics Reports* **489**, 89-159 (2010).
- [25] Xinzhi Liu, Lijie Hao, Yuntao Liu, Xiaobai Ma, Siqin Meng, Yuqing Li, Jianbo Gao, Hao Guo, Wenze Han, Kai Sun, Meimei Wu, Xiping Chen, Lei Xie, Frank Klose, Dongfeng Chen, *J. Magn. Magn. Mater.* **417**, 382 (2016).
- [26] E.F. Bertaut, Magnetism, edited by G. T. Rado and H. Suhl (Academic, New York, 1963), Vol. 3, pp.149.

- [27] Ganesh Kotnana and S. Narayana Jammalamadaka, *J. Appl. Phys.* **118**, 124101 (2015). Ganesh Kotnana, S. Narayana Jammalamadaka, *J. Magn. Magn. Mater.* **418**, 81-85 (2016).
- [28] Su Yuling, Zhang Jincang, Feng Zhenjie, Li Zijiong, Shen Yan, Cao Shixun, *J. Rare Earths* **29**, 1060 (2011).
- [29] S. Venugopalan, M. Dutta, A. K. Ramdas, and J. P. Remeika, *Phys. Rev. B* **31**, 1490 (1985).
- [30] I. B. Krynetskii, V. M. Matveev, *Phys. Solid State* **39**, 584 (1997).
- [31] Q. Zhang, J. H. Cho, B. Li, W. J. Hu, and Z. D. Zhang, *Appl. Phys. Lett.* **94**, 182501 (2009).
- [32] W. C. Koehler, E. O. Wollan, and M. K. Wilkinson, *Phys. Rev.* 118, 58 (1960).
- [33] Zhong Wei, Au Chak-Tong, and Du You-Wei, Chin. Phys. B 22, 057501 (2013).
- [34] T Chatterji, N Jalarvo, C M N Kumar, Y Xiao and Th Brückel, *J. Phys: Condens Matter* **25**, 286003 (2013).
- [35] M. Shao, S. Cao, S. Yuan, J. Shang, B. Kang, B. Lu, and J. Zhang, *Appl. Phys. Lett.* **100**, 222404 (2012).
- [36] M. C. Silva-Santana, C. A. daSilva, P. Barrozo, E. J. R. Plaza, L. de los Santos Valladares, N. O. Moreno, *J. Magn. Magn. Mater.* **401**, 612 (2016).
- [37] S. Yin, V. Sharma, A. Mcdannald, A. Reboredo, and M. Jain, RSC Adv 6, 9475 (2016).
- [38] Lingwei Li, Katsuhiko Nishimura, Wayne D. Hutchison, Zhenghong Qian, Dexuan Huo, *Appl. Phys. Lett.* **100**, 152403 (2012).
- [39] Lingwei Li, Takahiro Namiki, Dexuan Huo, Zhenghong Qian, and Katsuhiko Nishimura, *Appl. Phys. Lett.* **103**, 222405 (2013).

- [40] Lingwei Li, Jing Wang, Kunpeng Su, Dexuan Huo, Yang Qi, *J. Alloys Compd.* **658**, 500-504 (2016).
- [41] Venkatesh Chandragiri, Kartik K Iyer and E. V. Sampathkumaran, *Phys. Rev. B* **92**, 014407 (2015).
- [42] Mingjie Shao, Shixun Cao, Yabin Wang, Shujuan Yuan, Baojuan Kang, Jincang Zhang, *Solid State Commun.* **152**, 947 (2012).
- [43] O. Tegus, E. Brück, L. Zhang, Dagula, K. H. J. Buschow, F. R. de Boer, *Physica B* **319**, 174 (2002).
- [44] Arrott A, Phys. Rev. 108, 1394 (1957).
- [45] S. K. Banerjee, *Physics Letters* **12**, 16 (1964).
- [46] T. Yamaguchi, J. Phys. Chem. Solids 35, 479 (1974).

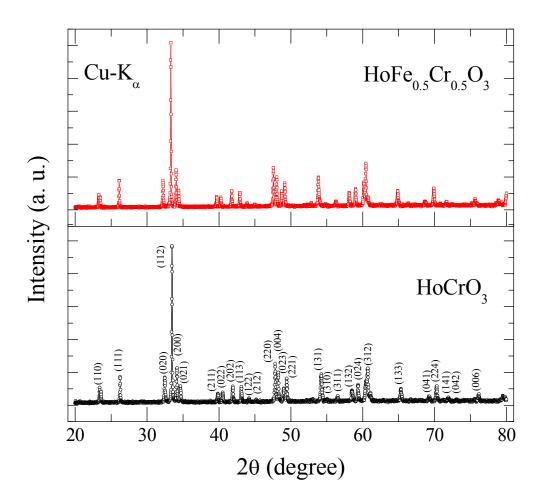


Fig.1: Powder x-ray diffraction patterns pertinent to HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ recorded at room temperature.

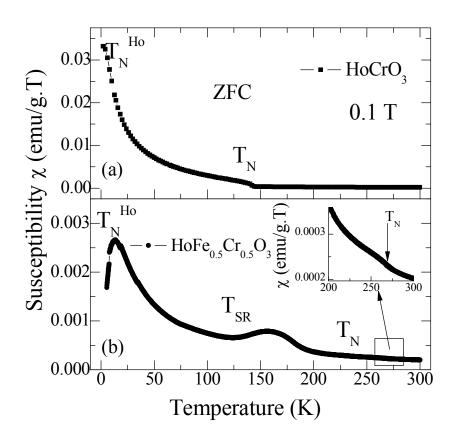


Fig. 2: Temperature (T) dependence of susceptibility (χ) pertinent to HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds measured at 0.1 T under zero field cooled condition (ZFC).

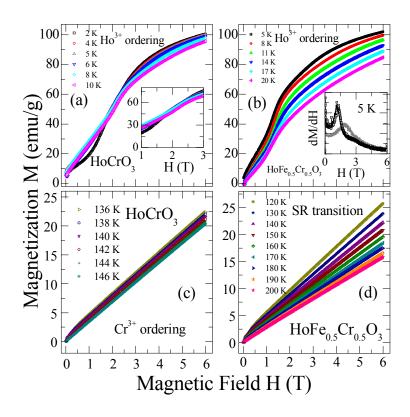


Fig. 3: (a - d) First quadrant magnetization isotherms near the transition temperature with an applied field up to 6 T for HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds. Magnetization crossover is shown as an inset of Fig.2 (a). Inset of Fig. 2(b) shows dM/dH vs. H for both the compounds corresponding to M vs. H measured at 5 K.

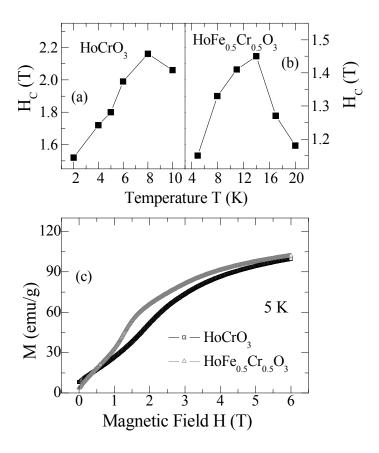


Fig. 4: (a) & (b) Temperature dependent variation of critical field (H_C) for $HoCrO_3$ and $HoCr_{0.5}Fe_{0.5}O_3$ compounds respectively (c) M vs. H graph measured at 5 K for both the compounds.

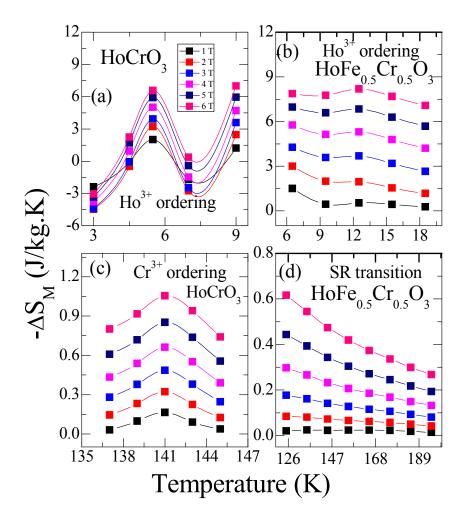


Fig. 5. (a – d) Temperature dependence of the magnetic entropy $-\Delta S_M$ obtained from magnetization M vs. H data for HoCrO₃ and HoCr_{0.5}Fe_{0.5}O₃ compounds.

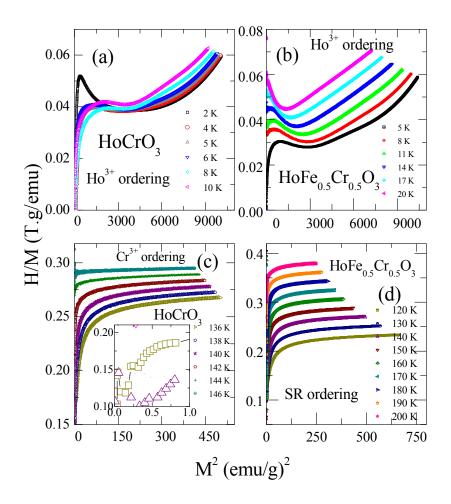


Fig. 6: Arrott plots corresponding to magnetization isotherms for $HoCrO_3$ and $HoCr_{0.5}Fe_{0.5}O_3$ compounds.