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Plasmonic Au NPs embedded Ytterbium-doped TiO₂ nanocomposites photoanodes for efficient indoor photovoltaic devices

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ABSTRACT

Herein, we present an all-inclusive investigation on the effect of plasmonic gold nanoparticles (Au NPs) embedded Yb-doped TiO₂ nanostructured photoanode for efficient photovoltaics devices. Optimally doped Yb (0.5 mol%)-TiO₂ nanoparticles (NPs) were synthesised by solid-state grinding. Their structural, morphological, optical, and surface properties were evaluated using XRD, Raman spectroscopy, SEM-EDX, UPS, and XPS. Results show that Yb-doped TiO₂ NPs form, reducing the rutile phase and shifting the fermi level. After demonstrating the enhanced properties of Yb-doped TiO₂ NPs, we introduced single anatase phase nanowire structures for comparison. Dielectric measurements show Yb-doped TiO₂ NW's higher conductivity. Plasmon-induced enhancement of light absorption is achieved by implanting Au NPs into Yb-doped TiO₂ NW for additional light harvesting. As a proof of concept, optimised Yb-doped TiO₂ photoanode DSSC were fabricated and tested under the AM1.5, 1 sun illumination exhibiting photoconversion efficiency (PCE) of 6 %. Devices employing Yb-doped TiO₂ NW: Au NPs photoanode exhibited J_{sc} to 18.8 % mA-cm⁻² and PCE% to 8 %. Under white LED (1000 LUX) exhibited PCE of 13.9 % for Au/Yb-TiO₂ NW photoanode. The electrochemical impedance spectroscopy (EIS) measurements on these devices established the Yb-doped TiO₂ NW: Au NPs photoanode based device exhibited the lowest charge transport resistance (R₂:86 Ω) and charge recombination resistance (R₃:20 Ω) which are consistent with the *J-V* characteristics suggesting the potential use of these photoanode for indoor photovoltaics (IPV) and in general any devices for allied energy and environmental areas of research.

1. Introduction

The advent of smart Internet of Things (IoT) devices in daily life have created a tremendous requirement to find alternative sustainable solutions to power these devices[1–4]. There is a great thrust to develop photovoltaic devices that work well in diffused light conditions, and the research area of indoor photovoltaics has received tremendous impetus recently[5–7]. Organic molecule-based photovoltaic devices, such as dye-sensitized solar cells[8,9], bulk heterojunction solar cells[10,11], and perovskite solar cells[12,13] have been tested for their utility under artificial indoor light. Rare-earth (RE³⁺) guests and semiconductor hosts are gaining attention for developing photoanodes with good optical and electrical properties. Notably, (RE) doped-TiO₂ is widely employed in photovoltaic and photocatalytic applications[14–17]. Optimal semiconductor materials are essential for photoconversion energy devices to maximize charge transfer efficiency, electrolyte diffusion, dye adsorption, and exciton generation[18,19]. TiO₂ semiconductor was widely

recognized as an effective mediator between dye molecules and electrodes [20,21]. However, the absorption capabilities were limited to the UV region because of the larger bandgap and defect states in the TiO₂ lattice, which inhibits the charge transport abilities. Therefore, TiO₂ has been modified with various dopants, including rare-earth metals, noble metal ions, transition metals, semiconductors, etc., to improve the optoelectronic properties of TiO₂ [22–27]. Among them, rare-earth metal dopants showed the greatest potential to tune the TiO₂ owing to their 4f-4f intra-band energy transitions, which led to improving the light-harvesting by upconversion/down conversion effects [28,29]. For example, Yb³⁺-TiO₂ (guest–host) has demonstrated an exceptional photoconversion energy efficiency (PCE%)[30,31]. With the doping effect of Yb³⁺, modifies the defect-sites in the TiO₂ lattice by substituting the Ti⁴⁺ with Yb³⁺. Thus, increases the number of oxygen vacancies, which facilitates into favorable band structure of TiO₂, thereby improving the charge transportation properties and low recombination rate at the photoanode-dye sensitizer interface[32]. Here in this, We

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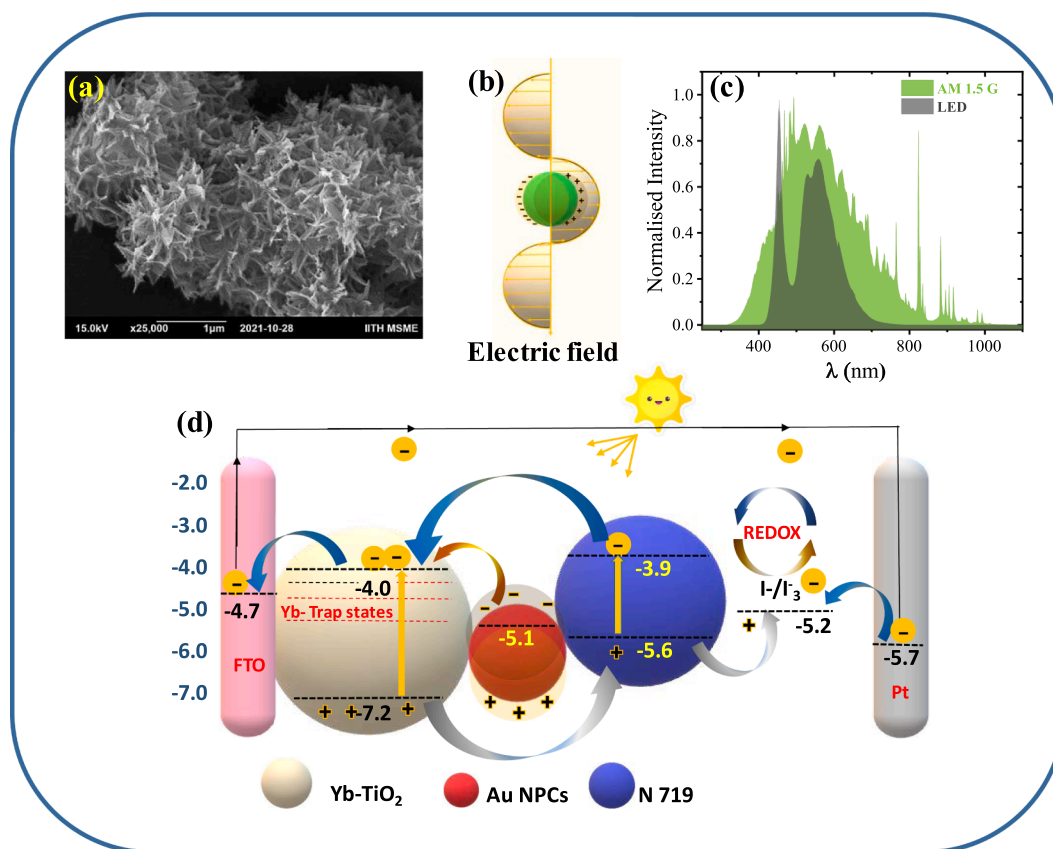
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used Yb (0.5 mol%)-TiO₂ to study the effect of Yb³⁺ on the intrinsic defects of TiO₂ lattice, with the intent of enhancing photovoltaic and photocatalytic capabilities[33–35].

Yb³⁺ is less beneficial for UV and visible light energy despite possessing the simplest energy structure. In these cases, integrating noble metals such as Au and Ag NPs into the Yb-TiO₂/N719 dye photoanode is a promising technique for enhancing the PCE% of DSSC[36]. It is known that the Au NPs induce a unique interaction between light and electrons on the surface of the metal nanoparticle. The phenomenon known as Local Surface Plasmonic Resonance (LSPR) occurs due to the oscillation of the electron cloud around the Au NPs in reaction to the applied electric field. Plasmonic nanoparticles exhibit substantial light absorption and scattering effects due to forming electric fields around them [37,38]. Indeed, the N719 dye molecules absorb light strongly in the visible spectrum, although their absorption deteriorates with increasing wavelength. Consequently, incorporating Plasmonic Au NPs is a significant step toward enhancing the overall light absorption efficiency, as gold nanoparticles have an extensive optical absorption effect in the visible region and strong optical interaction with the dye molecules[39]. Therefore, by utilizing the inclusive Plasmonic NPs, the resultant absorption coefficient of dye molecules can be improved [40]. Numerous groups have established the advancement of Plasmonic-induced DSSC [41–43]. In recent years, much research has been done on DSSCs containing multiple Plasmonic ions and several structural designs to improve photovoltaic performance[38,44–46]. In addition, the nanowire-structured TiO₂ can be replaced in DSSC as the nanowire can exhibit a greater surface-to-volume ratio, facilitating more dye adsorption, charge transport via the direct electron channel, and substantial light-harvesting via scattering effects [47–49]. The synthesized TiO₂ NW exhibits with single anatase phase, which is known to be providing the

higher photocatalytic activity[50]. Numerous groups have initiated the development of various nanostructured photoanodes for DSSC. Pawel et al. [51] demonstrated the comparative enhancement in the PCE% of DSSC with the various nanostructures of TiO₂, including nanoparticles, nanowires, and nanotubes of TiO₂. Furthermore, Huili et al. [23] noticed a 25 % increase in the PCE% of DSSC when a photoanode composed of noble metal (Au and Ag) was embedded in TiO₂ nanowires. M Kumari et al. [52] presented the Ag nanofibers embedded in TiO₂-based DSSC with 45 % enhancement over pristine TiO₂. Therefore, it is highlighted that the solar device efficiency can be improved by adding features such as Plasmonic materials and rare earth fluorescent doping[53–57]. For example, Peng Zhao et al. [58] showed a PCE% of 8.23 % for the Plasmonic DSSC employed with Au NPs decorated core cell structured NaYF₄: Yb, Er@SiO₂ photoanode. Whereas the NIR light is effectively trapped and converted to visible radiation by the up-conversion of rare-earth materials, this up-conversion is enhanced with the addition of Au NPs by superimposing the LSPR peak with the up-conversion emissions. Similarly, Parthian et al. [59] report an increase in PCE% of 21.3 % of DSSC using NaGdF₄: Yb, Er, Fe upconversion NPs and the Plasmonic silver NPs. In this study, we successfully used the synergistic effect of the LSPR effect of Au NPs and the rare-earth-fluorescent of Yb³⁺ to improve the performance of DSSCs. The summary of relevant literature has tabulated in Table S1. Hence, this study established a Plasmonic DSSC with a novel photoanode consisting of Au NPs embedded Yb³⁺-doped TiO₂ NWs as presented in Scheme 1. Multiple pathways of charge carrier generation and transport can be observed in the model, which explicitly improve the photovoltaic properties. Therefore, in this work, Au embedded Yb-doped TiO₂ nanowires (NW) structures photoanode based DSSC was constructed and achieved a PCE of 8 %, which is nearly 50 % higher than the pristine TiO₂ photoanode-based DSSC. In fact, no study



Scheme 1. Schematic representation of the plasmonic Au NPs induced DSSC with Yb-doped TiO₂ nanowire cluster structure photoanode under simulated Indoor (LED) light and Outdoor (AM1.5G) light conditions. (a) FESEM images of Yb-doped TiO₂ nanowire clusters; (b) schematic representation of Localised Surface Plasmonic Resonance (LSPR) effect; (c) spectral response of White LED light and solar light (Am 1.5G) spectra; (d) schematic representation of Plasmonic induced DSSC working.

has yet been published on the collaborative effort. As a result of DSSC's superior performance in low-light conditions, rather than direct sunlight [9,60,61], this new DSSC design paid attention to the indoor light photovoltaic technology. The research and development has been progressing on the indoor light based photovoltaic technology [30]. Indeed, the IPV system has been successfully used for various solar cell types, including organic photovoltaic (OPV) solar cells [62,63], Perovskite solar cells [12], and DSSC [7]. This work successfully performed the Plasmonic-induced DSSC under indoor light (white LED, 1000LUX). The Yb-doped TiO₂, Au embedded TiO₂, and nanowire clustered structured TiO₂ composite photoanode-based DSSC presents a higher PCE than the pristine TiO₂. In particular, the combined impact of Au NPs implanted Yb-doped TiO₂ NW-DSSC demonstrated a higher PCE (η) of 13.9 %, which is 66 % higher than undoped TiO₂ (η :8.6 %). Thus, the enhanced photovoltaic characteristics of under indoor (ambient) and outdoor (AM 1.5G) light conditions suggest a new path for designing a material combination for collaborative efforts to improve light harvesting. This discovery could pave the way for future research on low-light energy harvesting devices, as recycling indoor light energy could render energy-efficient electronics self-sufficient.

2. Experimental section

2.1. Chemicals

Titanium(IV) dioxide P25 (TiO₂, 99.96 %), Ytterbium (III) oxide (Yb₂O₃·5H₂O, 99.99 %), 4-*tert*-BUTYLPYRIDINE (C₉H₁₃N, 98 %), Lithium iodide (LiI, 99.9 %), Iodine (I₂, 99.99 %), 1,2-Dimethylimidazolium iodide (C₅H₉IN₂, 98 %), and Au NPs colloidal solution (20 nm, 100 %) were purchased from Sigma Aldrich (product no. 741965), N719 (Ruthenium complex, 95 %) dye molecules purchased from Great cell suppliers. Fluorine-doped Tin Oxide (8 Ω /square, LumTec, Taiwan) and Platinum (Pt) coated substrates were purchased from Solonix, Switzerland. Valeronitrile(CH₃(CH₂)₃CN, 99 %), Acetonitrile (CH₃CN, 99 %), acetylacetone (C₅H₈O₂, 99 %), DI water (18.2 M Ω cm⁻¹), Acetone (C₃H₆O, 99 %), Isopropanol (IPA, 99 %), and sodium hydroxide (NaOH, 98 %) pellets were purchased from the SRL supplier.

2.2. Yb-doped TiO₂ NPs synthesis

Yb_xTi_{1-x}O₂ synthesis and the photoanode preparation for DSSC were carried out as presented in the literature [64]. The optimal stoichiometry of Yb_{0.005}Ti_{0.995}O₂ was prepared through solid-state grinding for eight hours by mortar and pestle. The composition is then transferred to the furnace and annealed for two hours at 500 °C to reduce the composition's stress and hydrocarbon- ions contents. The final product is used for photoanode fabrication.

2.3. Yb- TiO₂ nanowire clusters via solvothermal method

Nanowire clusters were synthesized using the Solvothermal route reported in the literature [60] with few parameter changes. For Yb-doped TiO₂ NWS, solid-state hand-grinded Yb-doped TiO₂ is directly used as a precursor. In contrast, 4 gm of Yb-doped TiO₂ nanoparticles were added to a NaOH solution (10 M NaOH dissolved in 60 ml of DI water) at room temperature while stirring at 500 rpm. The reaction is arrested for more than one hour for homogenous mixing, and then the outcome precipitation is shifted into a 100-ml Teflon-lined autoclave. This autoclave was transferred to an oven and heated to 150 °C for a 24-hr in situ transformation into nanowires. The resultant Yb-doped TiO₂ nanowire clusters were taken out from the autoclave by fleshing the water and then washed with HCL (0.1 mol/L dispersing into 300 ml of DI water) for three hours. After that, the suspension is centrifuged to get a solid sample, and then it washes with DI water repeatedly till the pH value stabilizes to 7. The final composition is dried at 70 °C for 24 hr and then annealed at 500 °C for 30 min to remove the strain attained during

the synthesis process.

2.4. Characterization technique

X-ray diffraction spectroscopy (X'PertPRO with Cu K radiation of $\lambda=0.15406$ nm) was used to determine the structural characteristics. Additionally, lattice parameters and the cell size were estimated using the Rietveld refinement with FullProf software. Williamson-Hall (W-H) plots were used to estimate the grain size and strain values. Raman spectroscopy (Bruker Senterra dispersive Raman microscope spectrometer with 532 nm) was used to evaluate the molecular structure of the synthesized compositions. Surface analyses of the elemental composition were determined from XPS (Axis Model, Shimadzu) and FESEM-EDX (JEM 2100). Crystallite size and distribution were analyzed using TEM (JEOL F 200 COLD FEG-TEM), and further averaging the crystallite size by ImageJ software. The dielectric measurements were performed using the Wayne Kerr 6500B impedance analyzer over a frequency range of 1 Hz to 1 MHz on the pellet form compounds (pellet preparation S₃). UV-Visible spectroscopy (UV3092, LAB INDIA) was used to determine the optical absorption properties and calculate the bandgap values using Tauc-plot. UPS (Axis Model, Shimadzu) analysis determines the molecular orbital energy in the valance band. PL (FLS 1000, Edinburgh Instruments), and the Time-resolved Photoluminescence (TrPL) were performed using the Time correlated single photon counting (TCSPC) technique. Using a class AAA science tech solar simulator, the current density (J) versus potential (V) under (AM 1.5G, 1000 W/m²) was measured for the constructed DSSC, and a white LED (1000 LUX) light source was used for the IPV technology.

2.5. DSSC construction

FTO substrates were cleaned by sonicating with a series of solvents such as soap solution, DI water, acetone, Isopropanol for 15 min each. The cleaned substrates were masked with Scotch tape before film fabrication. The synthesized Yb-doped TiO₂ nanoparticles are converted into a paste by mixing with Triton-X and an acetylacetone-DI water binder (85 % – 15 %). The mixture was ground for 30 min using a mortar and pestle for a homogeneous paste. The doctor-blade technique was used to fabricate thin films, which were then subjected to a 15-min heat treatment at 70 °C and annealed at 500 °C for 30 min. Similarly, a second layer is deposited on top of it and then heated to 70 °C for 30 min, followed by annealing at 500 °C for 30 min. The films were further treated with TiCl₄ for 30 min at 70 °C to attain nonporous film, annealing at 500 °C for 30 min. Furthermore, 0.5 % M of gold nanoparticles were added to the Yb-TiO₂ paste during the pasting phase to evaluate the LSPR effect of gold nanoparticles on Yb-TiO₂ films. The films were then immersed in N719 dye for 12 hr, and the unadsorbed dye molecules were removed with acetonitrile devoid of water. The fabricated thin films served as photoanodes for the DSSC. TiO₂ deposited FTO substrate (working electrode) and platinum (counter electrode) are stacked using a surlyn polymer for DSSC construction. The electrolyte is composed of DMPII (1 M), I₂ (0.03 M), TBP(0.5 M), and LiI(0.05 M) dissolved in acetonitrile, which was injected into a device through a hole created on the back side of the counter electrode and sealed with thin glass sheet via heat treatment after electrolyte injection.

3. Results and discussion

3.1. Structural, optical and morphological properties of TiO₂ and Yb-TiO₂ nanoparticles

3.1.1. XRD

The phase purity and structural features of TiO₂ and Yb-doped TiO₂ were investigated using X-ray diffraction. The diffracted peaks for the TiO₂ are identified to be {101}, {103}, {004}, {112}, {200}, {105}, {211}, {204}, {116}, {220}, {215} at the position of 2θ : 25.33°,

37.1°, 37.99°, 38.63°, 48.25°, 54.16°, 55.30°, 62.86°, 69.09°, 70.49°, and 75.33°, respectively, indicating the anatase phase. Few low intense peaks are identified at 2 θ : 27.57°, 36.06°, 41.36°, 44.17°, 56.79°, and 62.88° that can be attributed to the lattice planes of {110}, {101}, {111}, {210}, {220}, and {002} respectively, which belongs to the rutile phase. Thus, the resultant XRD pattern of (P 25) reveals a dual phase of anatase and rutile phase, as depicted in Fig. 1(a). The resultant anatase and rutile phase were well matched with the standard data file (JCPDS 21-1272) and (JCPDS 21-1276), respectively. No additional peaks related to impurities or the second phase of Yb were detected in Yb-doped TiO₂. Significant phase matching with the XRD pattern of TiO₂ suggests that the Yb³⁺ dopant is effectively incorporated into the TiO₂ lattice[65]. Indeed, there is a negligible phase shift in the XRD pattern due to Yb-doping, as seen in Fig. 1(b), which may indicate the lattice deformation of TiO₂ caused by the substitution of Ti⁴⁺ with Yb³⁺. Furthermore, Rietveld refinement is accompanied on the XRD of Yb-doped TiO₂ to evidence the site replacement of Ti⁴⁺-ion with the Yb³⁺-ion. Fig. 1(c-d) demonstrates that the calculated XRD profile and observed data for undoped TiO₂ and Yb-doped TiO₂ were in good agreement. The calculated structural parameters are tabulated in Table 1. These structural modifications may lead to a reduction in strain and lattice constant[66]. Under Yb-doping into TiO₂, observed the reduction in the lattice parameter of “a” and “c”, where the reduced “c” parameter indicates the compression strain existence in the TiO₂ lattice [67]. The volume of the lattice cell is reduced to 135.88 (Å³) from 136.05 (Å³) for Yb-doped TiO₂. The reduced lattice cell of Yb-TiO₂ would signify the site substitution of Ti⁴⁺ with the higher ionic radius of Yb³⁺. When Yb-ions are incorporated into TiO₂, the production of Yb-Ti-O and Yb-O bonds is possible. This formation of new bonds could

Table 1

XRD profile structural parameters obtained from the Rietveld refinement on the XRD pattern of TiO₂ and Yb-doped TiO₂. “lattice parameters (a,b,c), and unit cell volume (V)”.

TiO ₂	Yb-doped TiO ₂
Anatase phase	Anatase phase
a = b = 3.783 Å	a = b = 3.782 Å
c = 9.507 Å	c = 9.500 Å
V = 136.05 (Å ³)	V = 135.88(Å ³)
Rutile phase	Rutile phase
a = b = 4.592 Å	a = b = 4.589 Å
c = 2.958 Å	c = 2.954 Å
V = 62.37(Å ³)	V = 60.20(Å ³)
Anatase (wt.%) = 92.97 % + Rutile (wt.%) = 7.03 %	Anatase (wt.%) = 97.25 % + Rutile (wt.%) = 2.75 %
R _p = 7.17	R _p = 12.5
R _{wp} = 9.92	R _{wp} = 14.1
χ ² = 3.6	χ ² = 2.85
Grain size = 39.9 nm	Grain size = 33.8 nm
Strain = 0.0034	Strain = 0.0020

stabilize the Ti-O bond, whereas the breakage of the Ti-O bond is required for grain growth/suppression. Hence, the phase transition and crystal growth are impeded[68]. Furthermore, when Yb³⁺ is substituted for Ti⁴⁺, there will be a lattice distortion of TiO₂ since Yb³⁺ has a larger ionic radius (1.0 Å) than Ti⁴⁺ (0.6 Å), and this may have contributed to the formation of interstitial titania, all these effects could lead to suppress the phase transition from anatase to rutile[69]. This suppression of phase transition under Yb-doping suggests grain inhibition since it is well known that the grain size of anatase phase TiO₂ is less than that of rutile phase TiO₂. Rietveld refinement analysis demonstrates the rutile

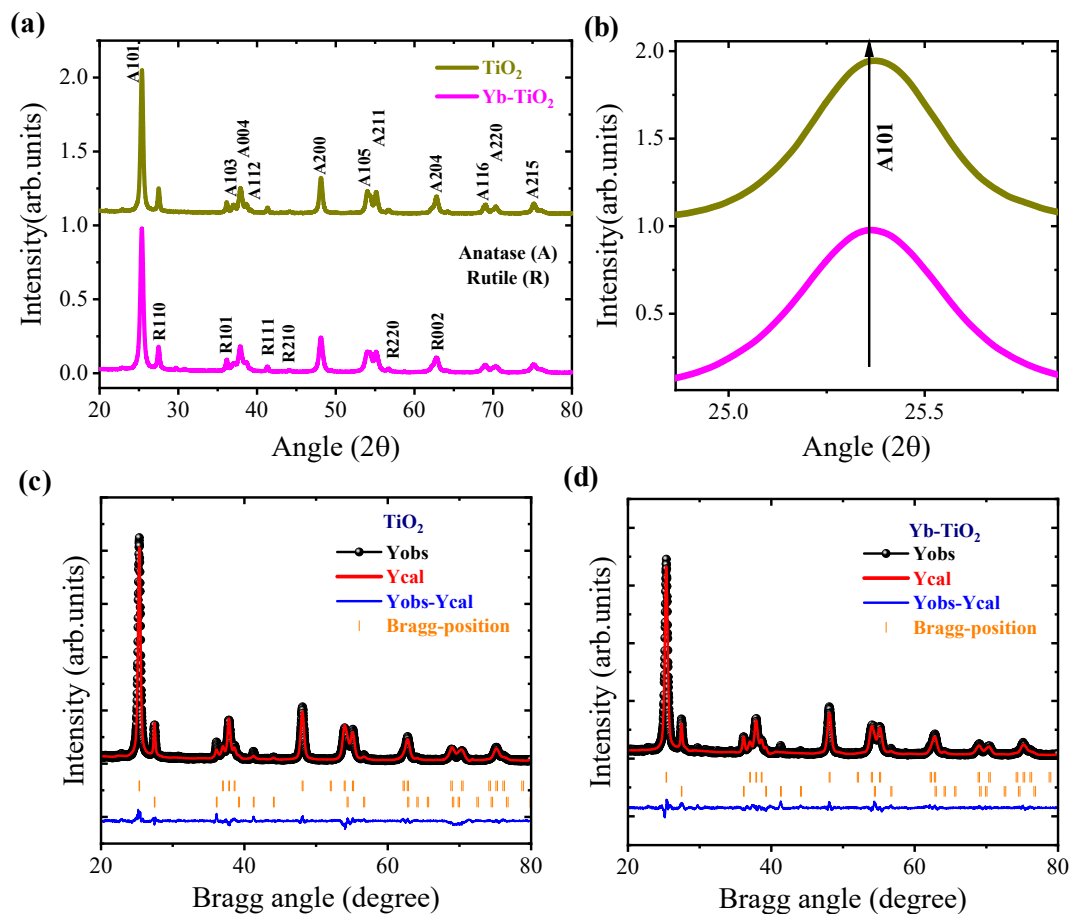


Fig. 1. (a) XRD comparison between Yb³⁺ doped TiO₂ and undoped TiO₂, (b) XRD peak (A101) shift representation, (c-d) Rietveld refinement quantitative analysis of XRD pattern of pristine TiO₂ and Yb-doped TiO₂.

phase decrease in the mixed-phase of TiO₂ under Yb-doping impact.

Williamson-Hall (W-H) plots distinguish the grain size and micro-strain for TiO₂ and Yb-doped TiO₂ under the following equation of $\beta\cos\theta = C\epsilon\sin\theta + (K\lambda/D)$ where ϵ is indicating the strain and D is the grain size (crystallites), the corresponding W-H plots presented in [supporting information S1](#). Yb-doped TiO₂ has smaller crystallites and less strain than TiO₂, as shown in [Table 1](#). In fact, small crystallites have an excellent surface-to-volume ratio, resulting in increased dye adsorption, which is essential for optoelectronic devices.

3.1.2. Raman spectroscopy

Raman spectroscopy was used to study the vibrational modes of Yb-doped TiO₂ and pristine TiO₂, as shown in [Fig. 2\(a\)](#). The Raman spectra of TiO₂ display Raman peaks at 142, 392, 515, and 635 cm⁻¹, which corresponds to the vibrational modes of symmetric stretching (E_g), and symmetric bending (B_{1g}), and anti-symmetric bending (A_{1g}), respectively. The Raman spectra of Yb-doped TiO₂ were observed to be similar to undoped TiO₂, indicating that the addition of doping elements does not influence the lattice structure of TiO₂ [34]. On the other hand, a slight phase shift was observed in the Yb-doped TiO₂ towards a longer wavelength, as shown in [Fig. 2\(b\)](#). there is quenching observed in the Raman peak intensity as a result of Yb doping, which is ascribed to the distorting of the Ti⁴⁺ centre and the formation of Ti³⁺ defects [70]. Thus, these results suggest that the Yb is effectively incorporated into the TiO₂ [71]. Indicating low crystallinity, which is consistent with the XRD analysis.

3.1.3. X-ray photoelectron spectroscopy (XPS)

The chemical analysis of the surface of Yb-doped TiO₂ and undoped TiO₂ samples was investigated using XPS. [Fig. 3\(a\)](#) depicts the XPS wide-scan spectrum, which distinguishes the available Ti (2P), O (1s), and C (1s) elements in TiO₂. In comparison, the Yb-doped TiO₂ displays a survey spectrum similar to the reported TiO₂. The Yb³⁺ elements enabled to detect over 182.1 eV region in the wide-scan spectrum indicate that Yb's formation on the surface of TiO₂ is significantly low. [Fig. 3\(b\)](#) depicts the high-resolution XPS spectrum of Ti 2P, representing spin-orbit doublets of Ti 2P_{1/2} at 464.5 eV and Ti 2P_{3/2} at 458.85 eV, which correspond to the Ti⁴⁺ oxidation state of TiO₂. Likewise, Yb-doped TiO₂ exhibits similar Ti 2P spectra to those of undoped TiO₂ with a slight peak shift towards lower binding energy. The shifted spin-orbit doublets were identified at 457.5 eV (Ti2P_{1/2}) and 463.5 eV (Ti2P_{3/2}), which may attribute to the Ti³⁺ oxidation state of TiO₂ [70].

As depicted in [Fig. 3\(c\)](#), the O 1s of both undoped TiO₂ and Yb-doped TiO₂ exhibit an asymmetric and widening spectrum. The O 1s orbital spectrum represents a strong peak at 530.05 eV, which corresponds to the Ti-O bond, and the other peak is identified at 531.8 eV, representing

the hydroxyl oxygen groups. The peaks were highlighted with the Gaussian peak fitting. In contrast, Yb-doped TiO₂ exhibited a phase shift in the O 1s orbital spectrum towards lower binding energy. The study indicates that the binding energy of Ti 2P orbital and O 1s orbital are reduced for Yb-doped TiO₂, indicating the formation of new chemical bonds between the Ti and Yb elements [72].

3.1.4. UV-Visible absorption spectroscopy

UV-vis absorption spectra were carried out to evaluate the optical absorption properties of undoped and Yb-doped TiO₂ compositions in Diffused Reflection Spectroscopy (DRS) mode. As illustrated in [Fig. 4\(a\)](#), inset), the optical edge of Yb-doped TiO₂ is redshifted, suggesting the extended absorption ability. Tauc plots with Kubelka Munk (K-M) function were employed to determine the band gap of the samples, as shown in [Fig. 4\(b\)](#). It was observed that the bandgap of Yb-doped TiO₂ is narrower than that of undoped TiO₂, probably due to trap-site alteration in TiO₂ by incorporating the Yb-intra-band states. Moreover, oxygen vacancies and Ti³⁺ oxidation states may develop when Yb³⁺ replaces Ti⁴⁺. Creating these additional states below the conduction band would cause the conduction band's edge to slip down. This engineered TiO₂ with a narrower bandgap improves absorption and charge injection via enhancing Ohmic contact. Formation of these additional states could improve photo-excited electron charge hopping. Thus, the charge recombination is expected to be low. UPS analysis was used to consolidate the Yb-doped TiO₂ band structure results.

3.1.5. Ultraviolet photoelectron spectroscopy (UPS)

UPS is carried out to investigate the influence of Yb dopant on the changes in the band structure of Yb-TiO₂. [Fig. 5\(a\)](#) illustrates the UPS of TiO₂ and Yb-doped TiO₂. Work function (ϕ) can be determined by deducing the secondary electron cut-off ($E_{cut-off}$) from the excitation energy ($h\nu$) of He I (21.22 eV). Accordingly, the ϕ can be expressed as $\phi = h\nu - (E_{cut-off})$ [73]. Whereas, $E_{cut-off}$ is obtained by extrapolating the higher binding energy edge intersection with UPS background as presented in [Fig. 5\(b\)](#). The work function of Yb-TiO₂ is calculated to be 4.11 eV, which is 0.2 eV larger than TiO₂ (3.91 eV), as shown in [Table 2](#). This improved work function may be attributable to the substitutional replacement of Ti⁴⁺ with Yb³⁺. Over the doping effect of Yb, there is a chance of forming the defect sites below the conduction band minima (CBM), which may belong to the impurity or second phase elements of Yb³⁺ energy states, thereby, fermi level shift occurred towards the valance band. Furthermore, ionization (E_i) energy and electron affinity (E_a) were estimated by following the equation of $E_i = \phi + E_{VBM}$, and $E_a = E_i - E_g$, whereas the valence band maxima (E_{VBM}) was obtained by extrapolating the low binding energy edge intersecting with the X-axis

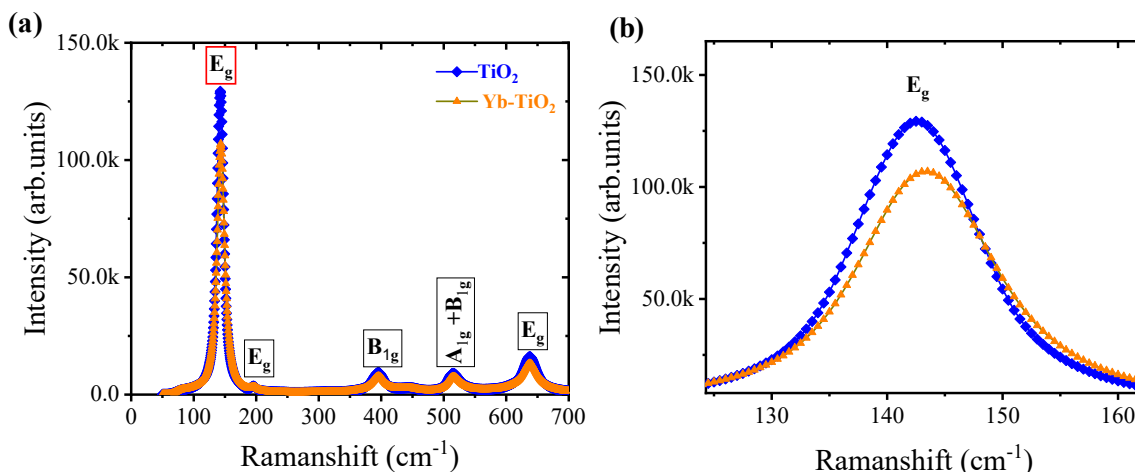


Fig. 2. (a) Raman spectra of Yb-doped TiO₂, and undoped TiO₂, (b) Raman peak (142 cm⁻¹) shift identification.

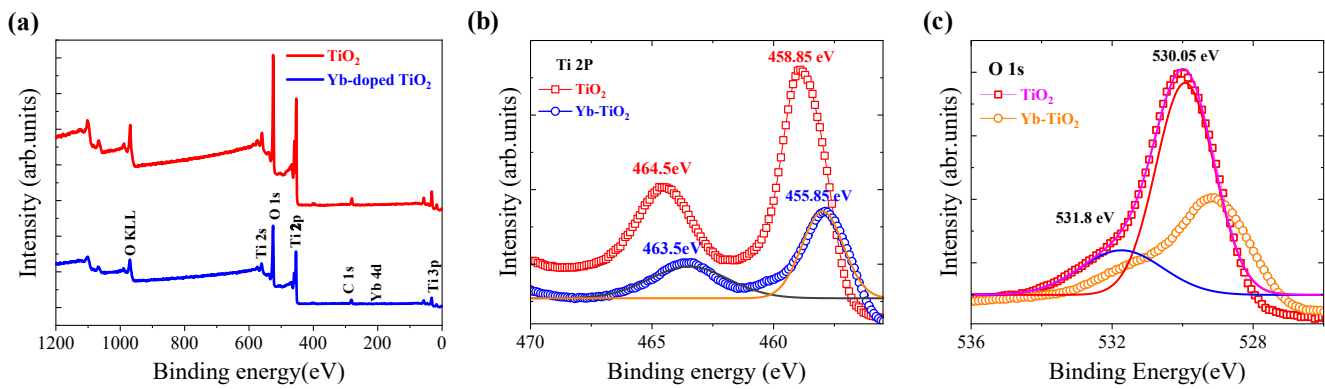


Fig. 3. XPS surface analysis. a) Wide-scan spectrum, (b-c) high-resolution XPS spectrum of Ti 2P and O 1s.

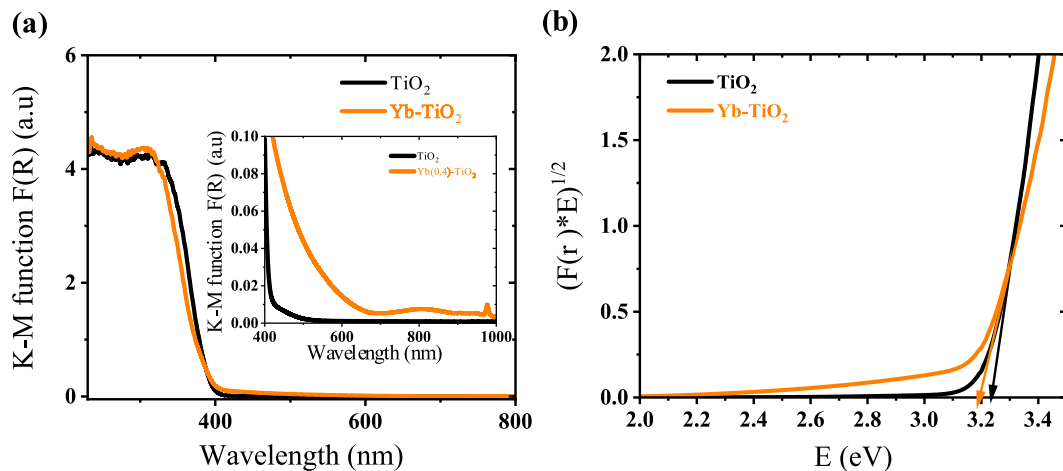


Fig. 4. (a) UV-visible absorption spectra in DRS mode (inset: absorption edge & Yb-related absorption peak at 980 nm) (b) Kubelka Munk plot for bandgap calculations of undoped TiO₂ and Yb-doped TiO₂.

as depicted in Fig. 5(c) [74] and E_g representing the bandgap of the sample as estimated by Tauc plots. Fig. 5(d) depicts the band-diagram for pristine TiO₂ and Yb-doped TiO₂ using all the calculated energy states. The band diagram reveals that the energy difference between the CBM and E_f is enlarged in the Yb-doped TiO₂, and the difference (CBM- E_f) is about 0.42 eV. This results can be relating to the formation of shallow/deep trap-states at below the conduction band. These trap-states are fundamentally crucial for charge transport and recombination effects. If the trap-site distribution range is near the CBM, the trap states are considered shallow traps. These trap states facilitate charge transport by enhancing the charge hopping mechanism [75]. Therefore, Yb-doped TiO₂ serves as an effective Electron Transport Layer (ETL) for the DSSC due to its enhanced charge transport capabilities. Indeed, the downshifted VBM implemented a built-in potential that enhanced charge migration at the interface. Then, Yb-doped TiO₂ is expected to have better charge transport and injection capabilities. The PL investigations add to evidence of improved charge transport and injection capabilities.

3.1.6. SEM and TEM images

Scanning Electron Microscopy (SEM) is performed on pristine TiO₂ and Yb-doped TiO₂ powder samples to investigate the morphological properties. TiO₂ exhibits a variety of particle sizes and shapes, as seen in Fig. 6(a-b). Spherical-shaped particles range in size from 10 nm to 60 nm, while nanorod-shaped particles range in length from 50 nm to 150 nm. Histogram plots display the particle size distribution, and the average particle size was 41 nm. The SEM image of Yb-doped TiO₂ shows smaller particles with an average size of 33.5 nm. It is usually

believed that bigger particles impose a high degree of crystallinity. Thus, Yb-doped TiO₂ has low crystallinity, and these values were consistent with the XRD plots. Furthermore, Transmission Electron Microscopy (TEM) was performed for better insight into the morphological properties, as shown in Fig. 6(c-d). As seen in the SEM images, TiO₂ crystallites came in various sizes and shapes. Crystallites range from 10 nm to 100 nm, averaging 37.4 nm. For Yb-doped TiO₂, the average crystallite size is 33.4 nm, and most of the crystallites were monitored in the range of 20 to 30 nm. Crystallite sizes calculated from TEM images agreed with those observed by XRD. In addition, they are measuring the selective area electron diffraction (SAED) pattern, as shown in Fig. 6(e-f). The displayed diffracted fringes of the Yb-doped TiO₂ and undoped TiO₂ samples illustrate the existence of poly-crystallinity. The phase purity and the structural existence can be distinguished using the diffracted fringes obtained from the SAED pattern. TiO₂ is identified with the lattice plans of {101}, {112}, {200}, and {211} based on the d-spacing values that have been determined from SAED pattern. The existing lattice plans indicate the presence of the anatase phase in the sample, and the additional {002} referring the rutile phase. Thus, the lattice fringes suggest that the TiO₂ contains anatase and rutile phases. The intensity of spots from Bragg reflection is reduced in the SAED pattern of Yb-doped TiO₂, suggesting the low crystallinity due to the doping effect. In summary, the obtained diffracted rings of TiO₂ and Yb-doped TiO₂ are well-matched with the XRD pattern.

3.2. Yb-doped TiO₂ nanowires

Fig. 7(a-b) presents the SEM images of Yb-doped TiO₂ NW at

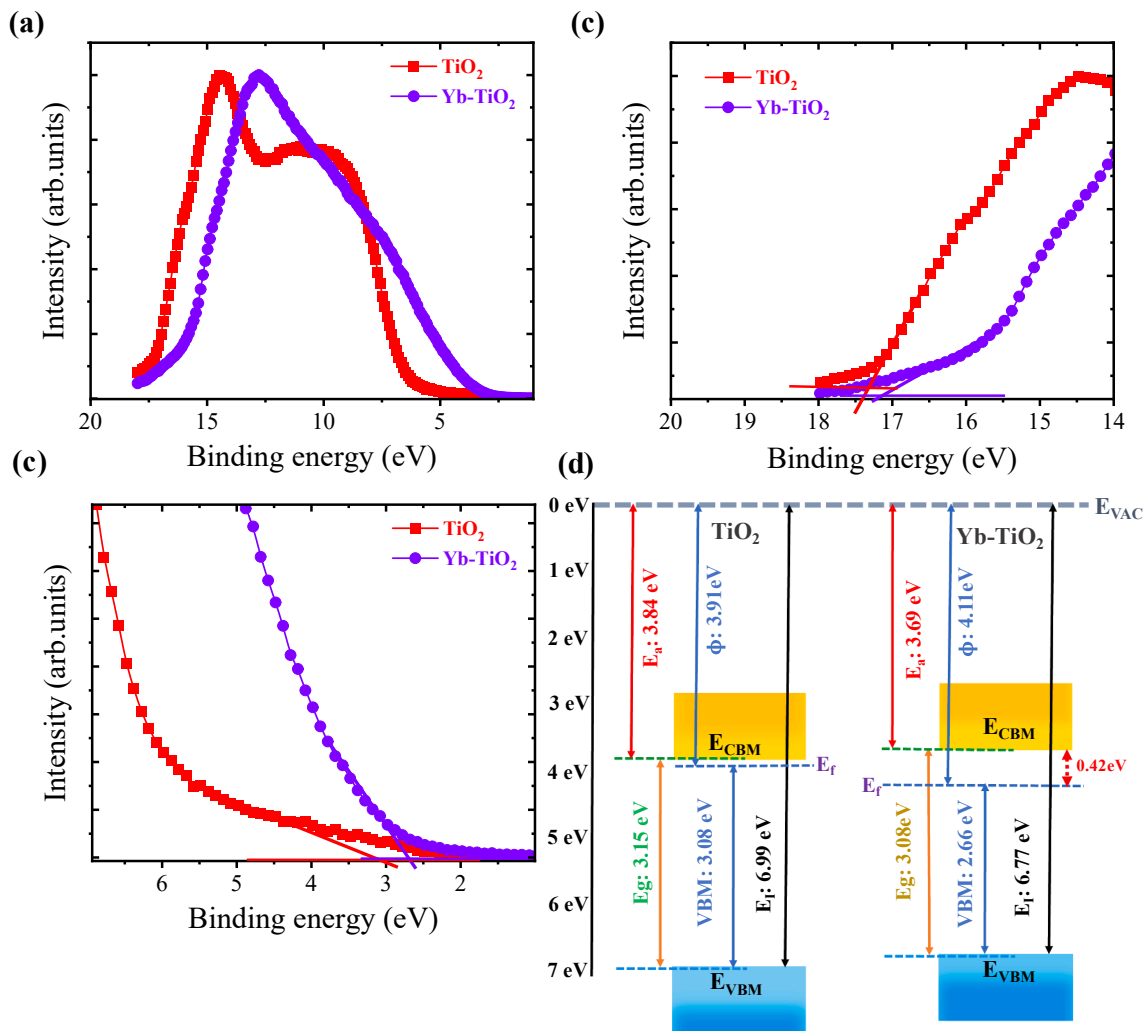


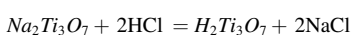
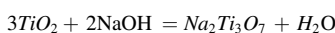
Fig. 5. (a) Ultraviolet photoelectron spectroscopy (UPS) analysis, (b) magnified view of secondary electron cut-off region (SECO), (c) valence band edge of undoped TiO₂, and Yb-doped TiO₂, and (d) band-diagram for pristine TiO₂ and Yb-doped TiO₂.

Table 2

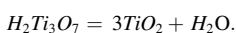
Calculated energy values of pristine TiO₂ and Yb-doped TiO₂ from UPS analysis.

sample	$E_{\text{cut-off}}$	E_{VBM}	ϕ	E_i	E_a	E_g	$E_f - \text{CBM}$
TiO ₂	17.31	3.08	3.91	44 nm	3.84	3.15	0.07
Yb-TiO ₂	17.11	2.66	4.11	41 nm	3.69	3.08	0.42

different magnifications. Here, a straightforward solvothermal approach converted the solid-state synthesized Yb-doped TiO₂ NPs into nanowires. Fig. 7(a) represents the SEM image, which resembles nano flowers or flakes. However, clusters of nanorods and nanowires can be observed in Fig. 7(b) when the image is zoomed in even more. The following equation depicts the formation of TiO₂ NWs during solvothermal treatment.[76].



After annealing the H₂Ti₃O₇ at 500 °C, all the hydrogen and additional oxygen ions are exempted, and then the TiO₂ anatase phase can be observed, as illustrated below.



The structural properties of undoped TiO₂ NWs and Yb-doped TiO₂

NWs were evaluated using the X-ray diffraction pattern shown in Fig. 7 (c). The XRD pattern of undoped TiO₂ NW exhibited a single anatase phase, which was well matched with the standard data file (JCPDS 21-1272). The XRD pattern of Yb-doped TiO₂ was identical to that of undoped TiO₂, showing that the solvothermal technique successfully introduced Yb into the TiO₂ lattice. In contrast to the TiO₂ NPs, the rutile phase has completely nucleated in the TiO₂ NW (supporting information S₂). These findings could be ascribed to the few atoms in the anatase lattice plane, which may produce nucleation sites for the rutile phase.[77]. In addition, the acid treatment that takes place during the solvothermal synthesis has the potential to nucleate the rutile phase of TiO₂ NW[78]. Furthermore, the elemental compositions of the materials were determined using the EDX spectrum, which is shown in Fig. 7(e-f). The EDX spectrum of TiO₂ reveals the elemental compositions of Ti and O, with no other impurity-related peaks found. Yb-doped TiO₂ yields a strong Ti peak and O peaks with few Yb³⁺-peaks. According to the quantitative analysis, Yb(III) has a stoichiometric doping percentage of (0.5 % mol), demonstrating the presence of Yb-peaks under ideal doping conditions.

3.3. Dielectric properties

The dielectric measurements were carried out for the undoped TiO₂, Yb-doped TiO₂ nanoparticles, and Yb-doped TiO₂ nanowire compositions over a frequency range from (1000 Hz to 10 MHz) at ambient

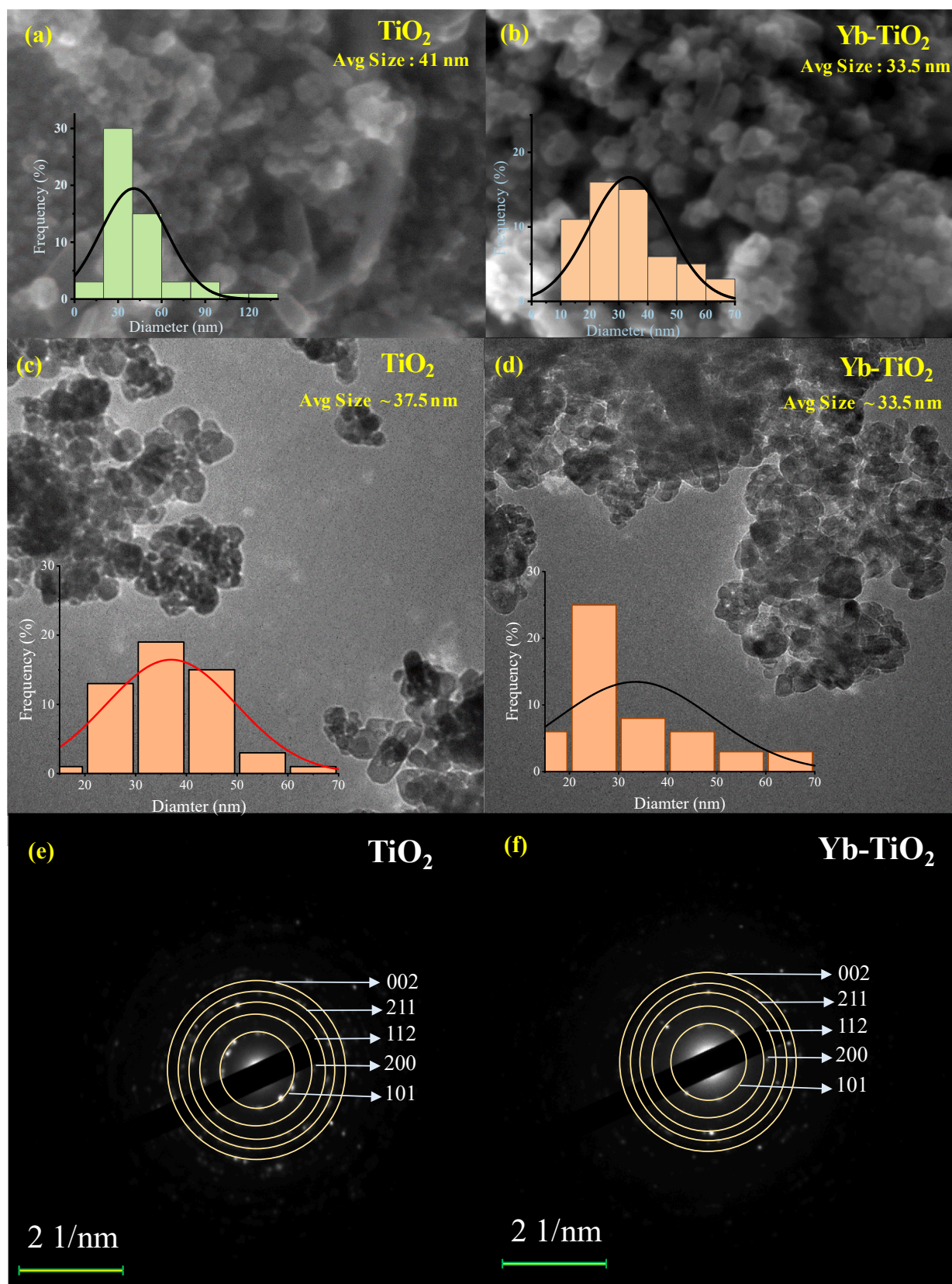


Fig. 6. (a-b) SEM images, (c-d) TEM images of pristine TiO₂ and Yb (0.5 M)-TiO₂ compounds (inset: histogram of the particle's diameter distribution), and (e-f) selective area electron diffraction (SAED) patterns of TiO₂ and Yb-doped TiO₂.

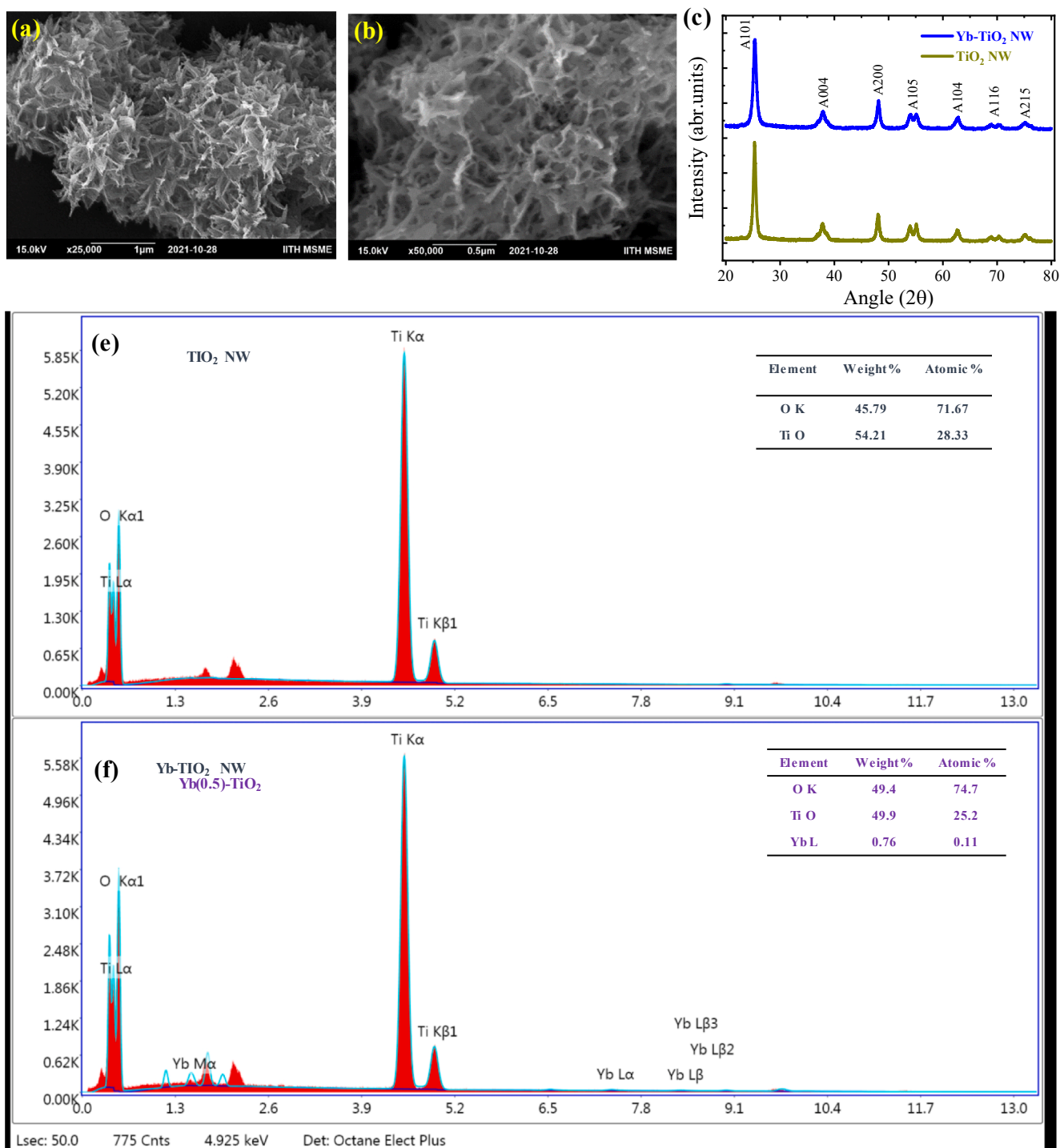


Fig. 7. (a-b) FESEM images of Yb-doped TiO₂ nanowire clusters (NWs) at different magnifications, Fig.(c) XRD pattern for pristine TiO₂ NWs and Yb-doped TiO₂ NWs synthesized via solvothermal route, and (e-f) Elemental compositions of TiO₂, and Yb-doped TiO₂ NWs.

conditions. The polarization of the dielectric material can be determined by the dielectric constant (ϵ'), estimated using the relation of $\epsilon' = Cd/\epsilon_0 A$. Where ϵ_0 indicates the permittivity of free space, C denotes the capacitance, d represents the thickness, and A denotes the pellet area. Fig. 8(a) displays the variation in the dielectric constant(ϵ') under a frequency range from 1000 Hz to 10 MHz. It has a higher dielectric constant at lower frequencies and gradually decreases as the frequency increases. The dielectric behavior of the material is determined based on the Maxwell-Wagner interfacial type of polarization. According to this model, the decrease in dielectric constant as frequency increases is

attributed to the formation of insulating grain boundaries around the grains. This phenomenon indicates that polarizability's effect is decreasing within the grain. Consequently, the dielectric properties diminish as frequency increases. As a result, it is clear that grain boundaries actively influence the (ϵ'), notably at low-frequency ranges. This higher ϵ' at lower frequency ranges can be attributed to space charge polarisation. On the other hand, grains can actively influence the ϵ' at midrange frequency, which is attributable to the dipole polarization relaxation. Since the grain size of the Yb-doped TiO₂ is smaller than undoped TiO₂ (evaluated from XRD), which introduces more grain

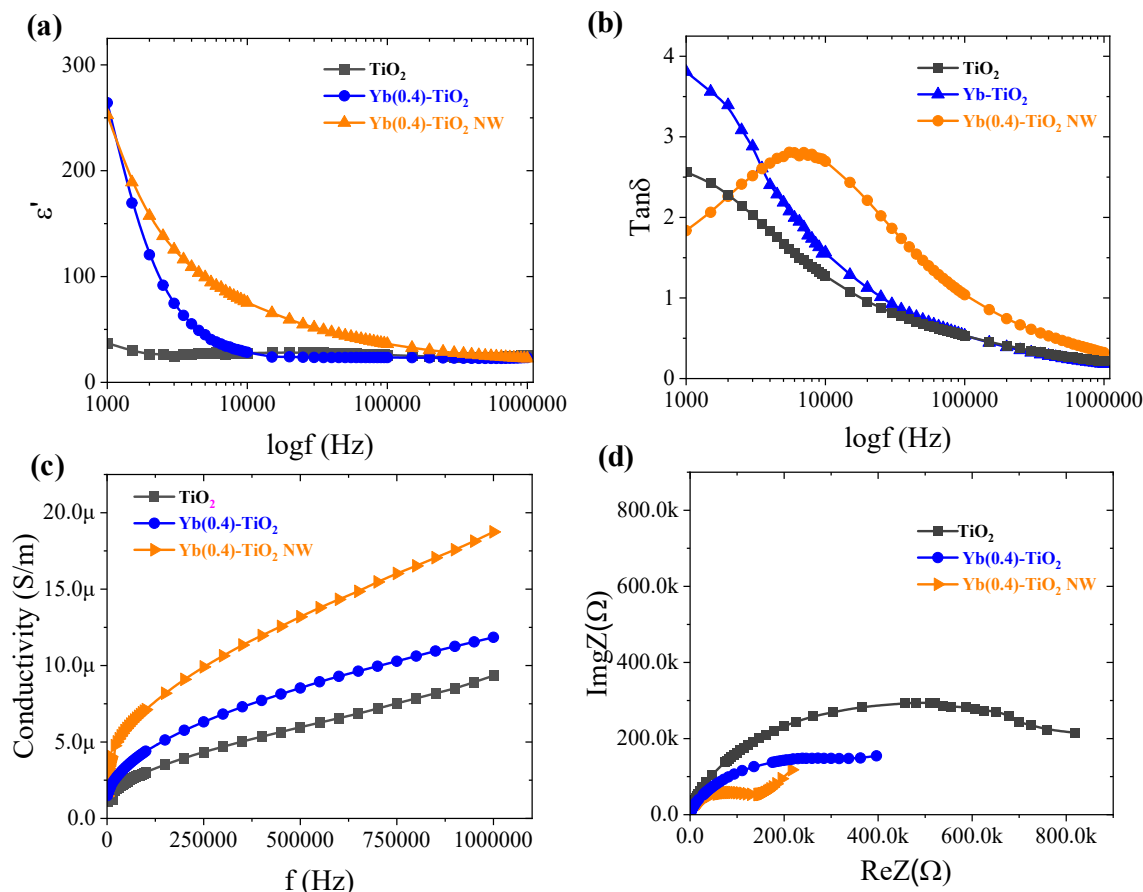


Fig. 8. (a) dielectric permittivity, (b) dielectric tangent loss, (c) electric conductivity, and (d) Nyquist plot of undoped TiO₂, Yb-doped TiO₂, and Yb-doped TiO₂ NW compositions under a frequency range of (1000 Hz to 10 MHz) at room temperature.

boundaries in the samples, thereby exhibiting the higher dielectric constant at a lower frequency regime. This study showed a sixfold enhancement in the dielectric constant for Yb-doped TiO₂ than pristine TiO₂ NPs at 1000 Hz. This enhanced dielectric constant is also attributed to the grain boundary defects and oxygen vacancies[79]. As mentioned earlier, oxygen vacancies are expected to form when Ti⁴⁺ replaces the Yb³⁺ oxidation state in Yb-doped TiO₂. Indeed, the dielectric constant can be attributable to the ability of electron hopping in the material[80]. The greater the number of electrons, the longer the hopping distance. Therefore, as a result of the oxygen vacancies present in Yb-doped TiO₂, the dielectric constant is substantially enhanced. Among all the samples, Yb-doped TiO₂ NW exhibits dielectric decay throughout the frequency range from 1KHz to 1 MHz, indicating the high polarise sensitivity.

Furthermore, the dielectric tangent loss was estimated using the equation of $\epsilon'' = \epsilon' \tan\theta$. As seen in Fig. 8(b), the dielectric loss ($\tan\theta$) exhibits a similar trend as ϵ' has demonstrated. Yb-doping enhanced dielectric loss at 10⁴, which can be attributed to the enhancement in the electric conductivity as the pace of charge conduction is responsible for the dielectric loss. Additionally, charge carrier mobility was evaluated using conductivity (σ) measurements, which account for electron transport abilities between grains and grain boundaries. Fig. 8(c) depicts the variations in the electric conductivity under the applied frequency. As frequency increases, conductivity increases, showing that the charges were mobilized in the presence of an electric field. Yb-doped TiO₂ NW exhibited a threefold increase in conductivity when compared to pristine TiO₂. Furthermore, the Nyquist plot (Z'' vs Z') was shown with the relevant frequency at ambient temperature. Fig. 8(d) illustrates the Nyquist plot with two semicircular arcs. The smaller semicircular arc at low frequencies is associated with grain boundary resistance effects, whereas the large semicircular arc at high frequencies is associated with

grain resistance. The smallest semicircular arcs were observed for Yb-doped TiO₂ NW, indicating increased charge conductivity at the grain boundaries.

In conclusion, the dielectric characteristics were comparable with the effective charge conductance of Yb-doped TiO₂ NPs/NWs compared to pure TiO₂. Solar cell applications benefit from materials with a high dielectric constant and electric conductivity.

3.4. Localized surface plasmonic resonance (LSPR) effect

The LSPR effect occurs when light strikes the metallic nanostructures. LSPR is a term that refers to the oscillation of an electron cloud surrounding metallic nanostructures. In detail, metallic NPs exposed to an electric field become polarised and behave as dipoles. The dipoles generate the restoring force within the nucleus and act opposite to the applied electric field. When the frequency of the applied electric field is resonant with the size of metallic nanoparticles, the dipoles exhibit the action of restoring force and oscillate with the applied electric field. Fig. 9(a) illustrates the metallic nanoparticles' LSPR effect. To establish the LSPR effect of Au NPs in Yb-TiO₂ composites, performed the UV-vis absorption measurements on Yb-doped TiO₂ thin-films after adding the Au NPs (0.5 mol %). The Au NPs coated thin film showing the surface Plasmon absorption band at 520 nm, as shown in Fig. 9(b). A comparable Plasmonic band can be observed for the Au NPs embedded Yb-TiO₂ thin-film compared to bare Yb-TiO₂. Therefore, it is emphasized that with the combined effect of dye absorption and the LSPR absorption band, more photons can be trapped, thereby improving the charge collection. Consequently, it was expected that the photovoltaic properties of Plasmon-induced DSSC would improve.

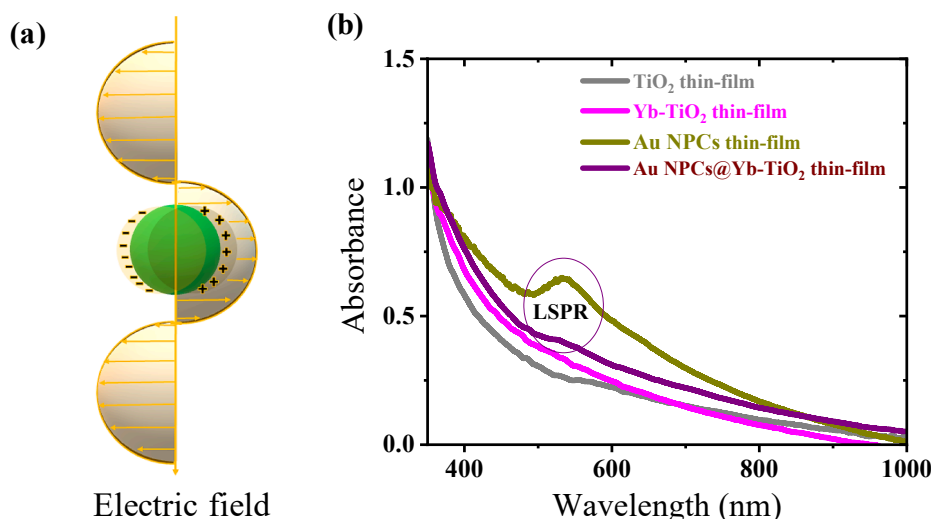


Fig. 9. (a) Schematic representation of Localised Surface Plasmonic Resonance (LSPR) effect, (b) UV-vis absorption spectrum of TiO_2 , and Yb-doped TiO_2 thin-film with LSPR absorption band.

3.5. Photoluminescence (PL) emission studies

Photoluminescence (PL) emission analyses were conducted to examine the charge carrier recombination/injection properties of ruthenium dye-loaded TiO_2 films. The ruthenium(N719) dye-loaded TiO_2 shows the PL emission at 720 nm under the excitation of 520 nm, as depicted in Fig. 10(a), which characterizes the recombination of excited e^-/h^+ pair in the dye molecules. The PL intensity of Yb-doped TiO_2 /dye is observed to be reduced as compared with TiO_2 /dye, this quenching effect indicating the lower recombination of excitons which improves the charge separation. As observed from the UPS analysis, the modified band structure of Yb-doped TiO_2 may improve Ohmic contact with dye molecules. This advantageous interface enhances the charge separation at the interface of dye molecules with Yb-doped TiO_2 [23]. Furthermore, with the addition of Au NPs into Yb-doped TiO_2 , the PL intensity reduced further, indicating that the charge recombination process has suppressed more under the influence of the Au NPs Plasmonic effect. Notably, the Schottky barrier at the interface between TiO_2 and Au NPs establishes the charge transfer path among them, thereby offering additional quenching in the PL emission. As estimated, the dye-sensitized TiO_2 NW showed lower intense PL emission than TiO_2 . Among them, Au-added Yb-doped TiO_2 NW reports a higher quench in the PL

emission. It is hypothesized that the combination of Yb-dopant, plasmonic Au NPs, and the nanostructured clusters effectively reduced the recombination process. To better understand, we have performed a Time-resolved Photoluminescence Lifetime (TrPL) upon the excitation of a 505 nm laser diode, as shown in Fig. 10(b). The ruthenium(N719) dye-loaded TiO_2 films present longer lifetime decays. The decay lifetime parameters were estimated using the Fluoracle software and tabulated in Table 3. The TrPL decay values are observed to become shorter for the engineered TiO_2 , suggesting the improved charge injection from dye molecules. Particularly, Yb- TiO_2 NW: Au photoanode reported effective charge injection abilities. Subsequently, Yb-doped TiO_2 NW: Au -based

Table 3

TrPL fitting parameters of N719 dye sensitised TiO_2 and doped TiO_2 thin-films.

Photoanode	τ_1 (ns)	τ_2 (ns)	τ_3 (ns)	χ^2
TiO_2	4	23	50	1
Yb- TiO_2	4.4	17.4	58	1
Yb- TiO_2 -Au	3.8	18	56	1
TiO_2 NW	2.8	15.3	48	1
TiO_2 NW-Au	2.7	15.2	46.3	1
Yb- TiO_2 NW-Au	3.2	17.7	54	1
Yb- TiO_2 NW-Au	2.6	14.3	46	1

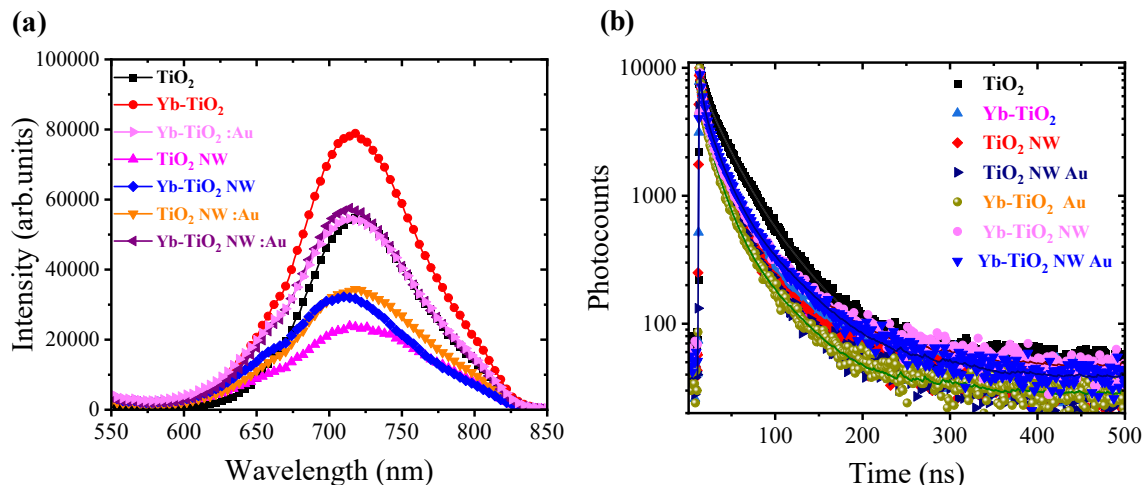


Fig. 10. (a-b) PL emissions and their electron decay lifetimes of N719 dye loaded TiO_2 , Yb-doped TiO_2 , and Au NPs added Yb-doped TiO_2 NWS thin-films.

DSSC exhibits exceptional enhancement in photovoltaic properties.

3.6. Photovoltaic characteristics of DSSC under AM1.5G, 1 sun illumination

Fig. 11(a) demonstrated the current density (J) vs voltage (V) parameters of DSSC. The associative photovoltaic parameters tabulated with the error analysis over four device sets are in Table. 4. Yb-doped TiO_2 -based DSSC exhibit higher PCE% than undoped TiO_2 . The enhanced PCE% can be attributed to the low recombination effect at the interface between dye molecules/ TiO_2 . As depicted in UPS, the engineered TiO_2 may have better Ohmic contact, and the formation of shallow traps improves the charge hopping. Moreover, incorporating Au NPs into the photoanode further increases PCE% with high current density (J). The improved (J) of Plasmon-induced DSSC can be attributed to the LSPR absorption band in conjunction with dye excitation. The DSSC with Au NPs added to Yb-doped TiO_2 exhibits a significantly higher PCE (η : 6.5%), which is about 20% greater than the test device's PCE (η : 5.3%), and the nanowire clusters TiO_2 photoanode-based DSSC increases the PCE% up to 5.9% with (J_{sc}) of $14.3 \text{ mA}\cdot\text{cm}^{-2}$. The DSSC with Plasmonic Au NPs embedded Yb-doped TiO_2 NWs-DSSC exhibits a higher current density of $18.8 \text{ mA}\cdot\text{cm}^{-2}$ with a PCE% of 8%. Specifically, the DSSC with a nanowire cluster structure yielded greater J_{sc} values due to the nanowires' greater surface-to-volume ratio, which facilitated enhanced dye molecule adsorption. Fig. 11(b) depicts the IPCE spectral response of the undoped and Yb-doped TiO_2 photoanode-based DSSC. The IPCE spectrum of the Yb-doped TiO_2 /N719 photoanode-based DSSC is improved due to a favorable trapping system in the Yb-doped TiO_2 lattice, allowing for less recombination. Au NPs

Table 4

photovoltaic parameters of the Au NPs induced Plasmonic DSSC with nanowire cluster structured photoanode, tested under one sun (AM 1.5G) light illumination.

Photoanode	V_{oc} (V)	J_{sc} ($\text{mA}\cdot\text{cm}^{-2}$)	FF	η (%)
TiO_2	0.7 ± 0.03	12.16 ± 0.1	62 ± 1.2	5.3 ± 0.2
TiO_2 :Au NPs	0.71 ± 0.05	14.3 ± 1.1	55 ± 2.3	5.6 ± 0.03
Yb- TiO_2	0.71 ± 0.1	13.7 ± 0.05	61 ± 0.6	6.1 ± 0.1
Yb- TiO_2 : Au NPs	0.7 ± 0.04	14.4 ± 0.124	63 ± 3.9	6.55 ± 0.02
TiO_2 NW	0.71 ± 0.06	14 ± 0.9	59 ± 2.5	5.9 ± 0.05
TiO_2 NW: Au NPs	0.7 ± 0.05	15.7 ± 0.06	59 ± 3.2	6.38 ± 0.12
Yb- TiO_2 NW	0.7 ± 0	17 ± 0.2	63 ± 1.3	7.5 ± 0.09
Yb- TiO_2 : NW Au NPs	0.72 ± 0.07	18.8 ± 0.04	59 ± 4	8 ± 0.21

Plasmon-induced DSSC had a more significant IPCE percentage, which refers to an increase in the optical density surrounding the metal surfaces in SPR-induced DSSC. As a result, more photons have been harvested [58]. The relative IPCE spectra ($(\Delta\text{IPCE}/\text{IPCE}) * 100\%$) are calculated to estimate the improvement in photo spectral response with the respective test device (TiO_2 -based DSSC) under the influence of Yb-doping and Au NPs. Fig. 11(c) illustrates that the relative IPCE is distributed over 300nm to 800 nm range. Particularly, the intensity of the IPCE spectrum from 400 nm to 600 nm is increased for Au NPs embedded TiO_2 and Au-embedded Yb-doped TiO_2 . The improvement in the IPCE spectrum is attributable to the LSPR effect of Au NPs at 520 nm (Fig. 7) [81,82]. With the doping effect of Yb^{3+} , the extended relative IPCE spectra in the IR region demonstrate enhanced charge collection. Table S1 (in Supplementary information) presents a comparative performance of Au NPs Plasmon induced Yb-doped TiO_2 @ Au NPs/N719

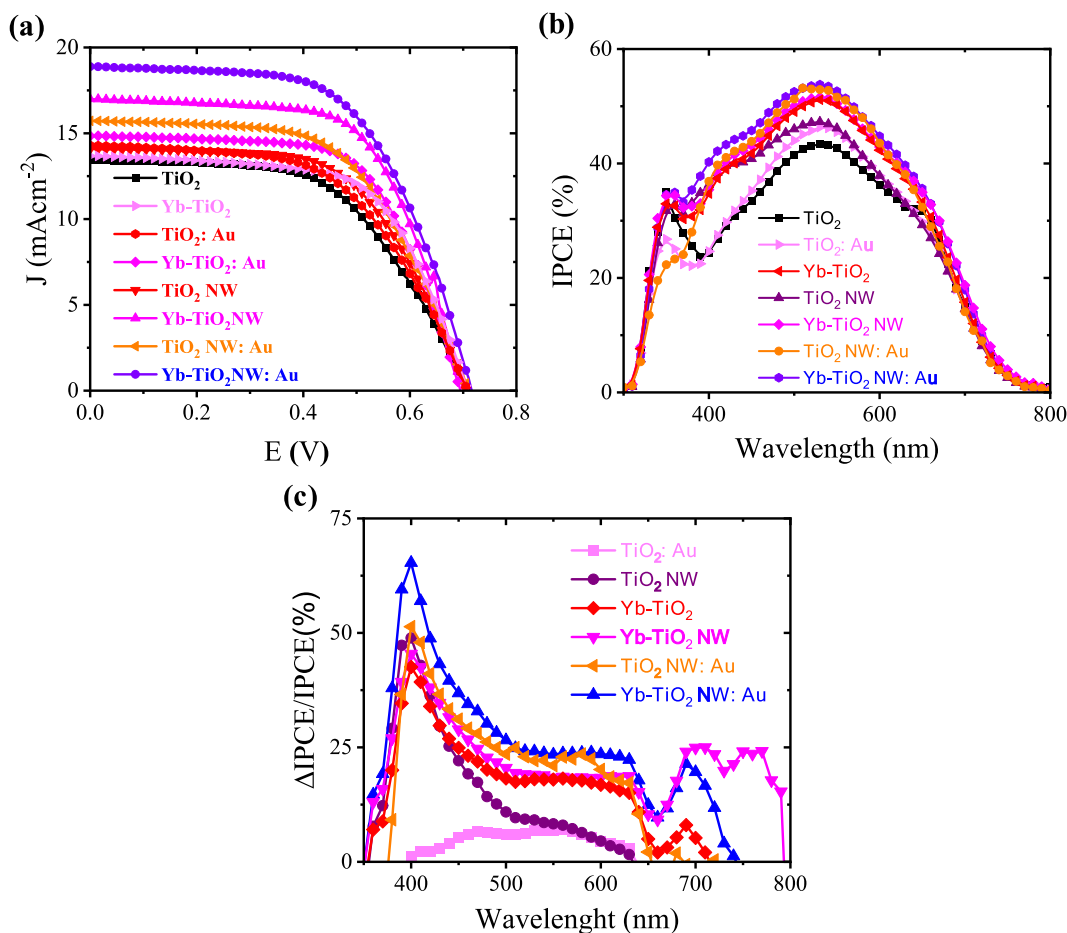


Fig. 11. (a). J .vs V plot, (b) IPCE spectral response, and (c) Relative changes in the IPCE spectral response of DSSC with nanostructured TiO_2 , Yb-doped TiO_2 , and Au NPs embedded Yb-doped TiO_2 photoanodes.

dye photoanode based DSSC vis-à-vis to various similar system from literature.

3.7. Indoor light photovoltaic (IPV) applications

The DSSC was then tested under indoor lighting conditions, while its optimal photovoltaic characteristics were monitored under (outdoor) AM 1.5G (1sun, 100 mW/cm²) illumination. The light intensity is set for indoor light conditions to 1000LUX(327 μW/cm²) using a white LED light source[7]. The ruthenium dye-sensitized solar cells have presented better results than thin-film and Si-based solar cells in indoor light conditions [83,84]. Fig. 12(a) presents the ruthenium(N719) dye absorption spectrum that overlaps with the spectral response of AM 1.5G and the white LED source. It was observed that the spectral response of LED was completely merged with the absorption spectrum of N719 dye molecules. Therefore, the DSSC could effectively generate more no excitations under LED illumination, thereby improving the photovoltaic properties.

Fig. 12(b) depicts the J-V characteristics of DSSC under White LED illumination with different nanostructured TiO₂ and Yb-TiO₂ photoanodes. Whereas Yb-doped TiO₂ employed DSSC reported with higher efficiency than pristine TiO₂ photoanodes, the corresponding parameters are tabulated in Table 5. With the addition of Plasmonic Au NPs, the DSSC was extended further into a higher PCE% of 11.2 %. Consequently, the Au/Yb-TiO₂ NW/dye photoanode increased the PCE% by more than 65 percent, and greatly enhancing photocurrent density compared to the test device.

3.8. EIS-Nyquist plots of DSSC

Electrochemical impedance spectroscopy (EIS) studies were used to characterize the interfacial charge transport resistance of the device. The EIS-Nyquist plot was performed under a frequency range of 1 MHz to 0.1 Hz, and an open-circuit voltage of 0.7 V under illumination by a White LED light source. Fig. 13 shows the Nyquist plot of DSSC, representing three depressed semicircles. The pertinent equivalent circuit was used to fit the Nyquist plot and extract the EIS parameters, as shown in Fig. 11(inset). In contrast, the first semicircle at high frequency is attributed to charge transfer resistance(R₂) at the counter electrode/electrolyte interface. The mid-range frequency semicircle represents the charge recombination resistance (R₃) at the TiO₂ photoanode/electrolyte interface. The third semicircle at a lower frequency is related to the Warburg diffusion resistance of electrolytes [85,86]. The distance between the origin and the Nyquist plot starting point on the left side indicates the series resistance (R₁). Table 6 lists the EIS parameters that were computed from the Nyquist plot. In contrast, the series resistance

Table 5

photovoltaic parameters of the constructed DSSC under white LED (1000 LUX, 327 μW/cm²) illumination.

Photoanode	V _{oc} (V)	J _{sc} (μA-cm ⁻²)	FF	η(%)
TiO ₂	0.56 ± 0.01	135 ± 2.1	37 ± 0.2	8.6 ± 0.1
Yb-TiO ₂	0.58 ± 0.002	129 ± 3.4	48 ± 0.4	11.15 ± 0.3
Yb-TiO ₂ : Au NPs	0.59 ± 0.01	155 ± 1.6	47 ± 3.7	13.3 ± 0.12
TiO ₂ NW	0.565 ± 0.03	144 ± 5.9	40 ± 3.3	9.82 ± 0.5
Yb-TiO ₂ : NW	0.59 ± 0.01	148 ± 4.0	42 ± 2.1	11.4 ± 0.5
Yb-TiO ₂ : NW Au NPs	0.59 ± 0	153 ± 3.2	50 ± 0.1	13.95 ± 0.3

of (R₁) is identical across all devices, indicating that the components used to assemble the devices were optimized. R₂ and R₃ were exceptionally high in pristine TiO₂ (R₂:425 Ω and R₃:701 Ω, respectively). Yb-dopant and Au NPs influenced DSSC showed the lower interfacial resistances. Consequently, the charge transport is improved, which is crucial for enhanced photovoltaic characteristics. Specifically, Au NPs embedded Yb-doped TiO₂ NW exhibited the lowest charge transport resistance (R₂:86 Ω) and charge recombination resistance (R₃:20 Ω). These results were consistent with the J-V characteristics.

3.9. LSPR-induced DSSC

Fig. 14 displays the schematic representation of Plasmon-induced Yb-doped TiO₂ photoanode-based DSSC and the photoelectron generation/transport prospects. In this device, the dye molecules generate the electron/hole (e⁻/h⁺) pairs by absorbing the photons. Simultaneously, the Au NPs embedded in TiO₂ generate the (e⁻/h⁺) by creating a strong electric field around the Au NPs. This electric field is attributable to near-field effects with a penetration depth of up to 10 nm [87]. In contrast, due to the interaction with the Au NPs, the fermi level of TiO₂ is lowered to match the fermi energy level of Au NPs. And this down-shifted fermi level is able to enhance the charge injection efficiency from the dye molecules to the TiO₂ conduction band. Similarly, Under the influence of the Yb-dopant, the fermi level shifted down due to the formation of trap sites, as demonstrated in the UPS analysis, and these modified energy levels and formation of trap states are beneficial for charge transportation and injection at the interface with dye molecules. Thus, the high PCE% of Plasmon-induced DSSC is attributable to the combined LSPR absorption/scattering effect and defect-tuned energy band of Yb-doped TiO₂.

4. Conclusion

In conclusion, we have demonstrated the ensuing way to improve the PCE% by tuning the band structure of TiO₂ under the influence of Yb-

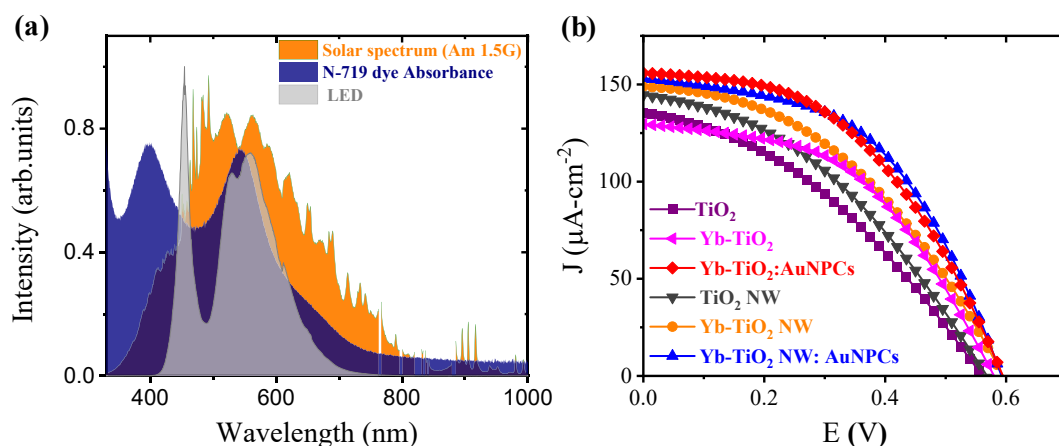


Fig. 12. (a) Spectral overlapping of the absorption spectrum of N719 dye molecules with the irradiance and emission spectrum of solar light (Am 1.5G), and White LED light, (b) J - V characteristic plots under Indoor light (White LED, 327 μW/cm²) at ambient condition.

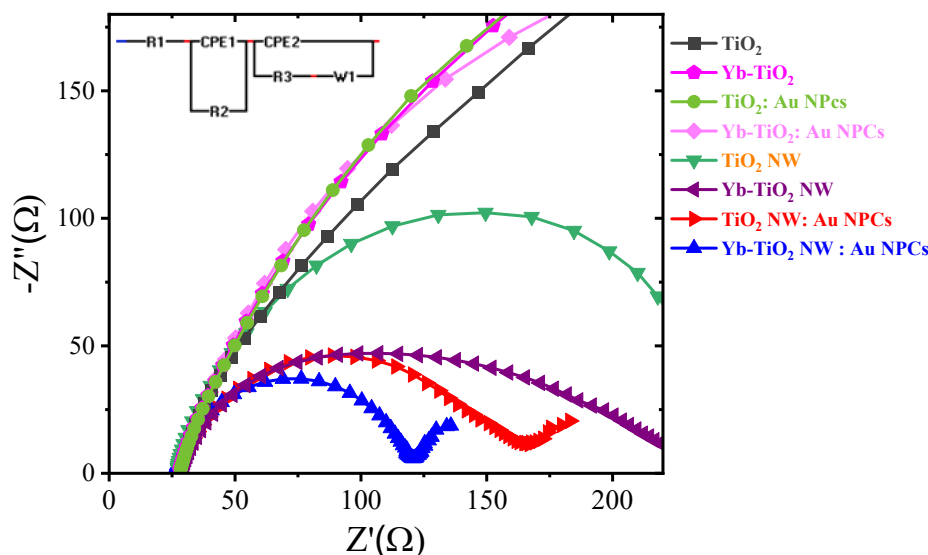


Fig. 13. Nyquist plots of DSSC based on various nanostructured TiO₂, Yb-doped TiO₂, and Au NPs embedded Yb-doped TiO₂ photoanodes.

Table 6

EIS parameters of TiO₂, Yb-doped TiO₂, and Au NPs added Yb-doped TiO₂ photoanodes applied DSSC under white LED (30 mW/cm²) illumination.

Photoanode	R ₁ (Ω)	R ₂ (Ω)	R ₃ (Ω)
TiO ₂	24	425	701
Yb-TiO ₂	30	129	606
TiO ₂ : Au NPs	21	556	333
Yb-TiO ₂ : Au NPs	27	418	297
TiO ₂ NW	26	40	150
Yb-TiO ₂ : NW	29	74	109
TiO ₂ : NW Au NPs	28	104	29
Yb-TiO ₂ : NW Au NPs	26	86	20

dopant, plasmonic Au NPs. Plasmonic-induced DSSC with Au NPs added Yb-doped TiO₂ NW have constructed and executed with tremendous PCE%. Here, we have synthesized Yb-doped TiO₂ nanowire clusters and utilized them for photoanode fabrication. With the XRD, Raman, XPS, FESEM-EDX, and TEM analysis, we have confirmed the successful doping effect of Yb into TiO₂ lattice with improved properties. The Rietveld refinement analysis has shown the reduced lattice parameters

of 'a' and 'c' for Yb-doped TiO₂, which evidence the site substitution of Ti⁴⁺ with Yb³⁺. Crystallite size & its distribution, variation in the phase transformations, and elemental composition & the formation of defect sites in the TiO₂ lattice were demonstrated in this work. Furthermore, dielectric characteristics illustrate the improved polarizability and electrical conductivity with the effect of Yb-doping. Likewise, the absorption peaks were found to be redshifted and identified with the reduced bandgap for Yb-doped TiO₂. By integrating the UV-vis and UPS analysis, the shallow trap states have been observed below the conduction band of TiO₂ by doping with Yb-intermediate energy states, which effectively improved the charge hopping mechanism and enhanced the photovoltaic properties of DSSC. PL quenching effects and short lifetime decay values ascribe efficient charge transportation/injection properties. Au NPs added to Yb-doped TiO₂ showed enhanced charge injection capabilities with dye molecules. Upon exposure to light, the electrons get excited and jump at the semiconductor/metallic nanoparticle interface, bringing the Fermi levels of TiO₂ and Au NPs into alignment. The downshifted fermi level of the Au-embedded Yb-doped TiO₂ enables the successful injection of charge from the dye molecules, as indicated by the PL quenching effect. As a result, the photovoltaic

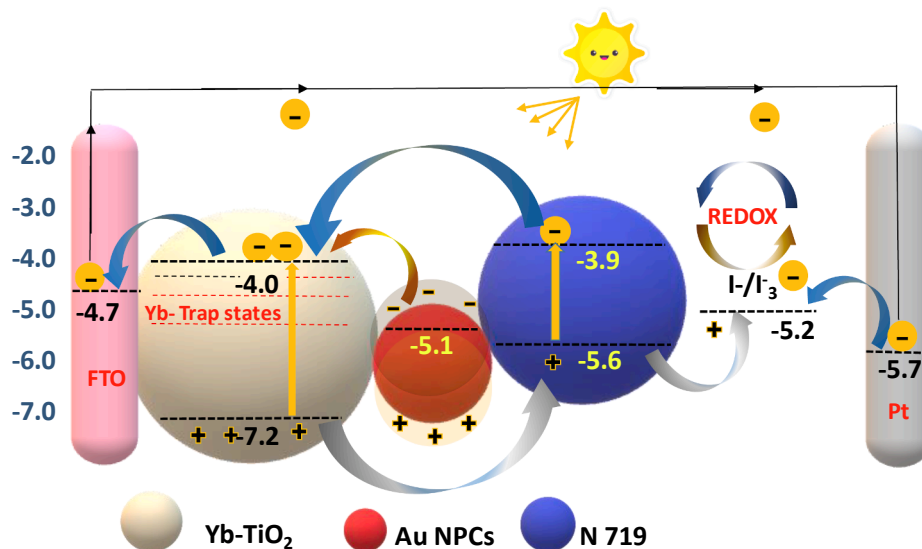


Fig. 14. Schematic representation of Plasmonic induced DSSC working.

properties were essentially enhanced by the influence of Yb-dopant and Au NPs. In contrast, Au NPs embedded Yb-doped TiO₂ photoanode-based DSSC exhibits a 20 % higher PCE% than the reference device (5.3 %). Additionally, when nanowire cluster structures are used in place of nanoparticles in the photoanode design, the PCE increases by more than 40 % compared to the test device. The Au-added Yb-TiO₂ NW performed well for indoor light conditions, with a PCE enhancement of roughly 52 % above undoped TiO₂ (η : 8.6 %). The enhancement in the light-harvesting abilities of DSSC was consistent with the IPCE spectral response. The EIS-Nyquist showed the decreased interfacial charge transport resistance for Plasmonic DSSC employed with Au-embedded Yb-TiO₂ nanowire clusters. As a result, the findings demonstrate a unique approach for enhancing the absorption capacity of TiO₂ by including rare-earth + 3 oxidized Yb(III) and Plasmonic Au NPs to harvest more solar energy.

CRedit authorship contribution statement

Venkata Seshaiha Katta: Data curation, Formal analysis, Investigation, Validation, Writing – review & editing, Writing – original draft. **V.R. Chappidi:** Formal analysis, Investigation, Validation, Methodology, Writing – original draft, Writing – review & editing. **Sai Santhosh Kumar Raavi:** Conceptualization, Resources, Funding acquisition, Supervision, Formal analysis, Investigation, Validation, Methodology, Writing – original draft, Writing – review & editing.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

Data availability

Data will be made available on request.

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Appendix A. Supplementary material

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.apsusc.2022.155728>.

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